

Weak Gravitational Lensing as a Tool for Cosmology

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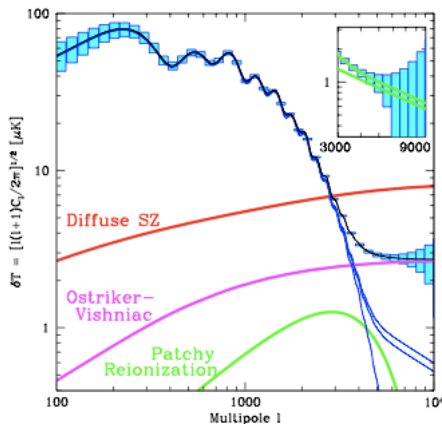
Information encoded in the CMB

At large angular scales, $\theta > 5'$ or $\ell < 2500$:

- primary anisotropies: small linear perturbations in the primordial plasma's density and velocity.

At small angular scales:

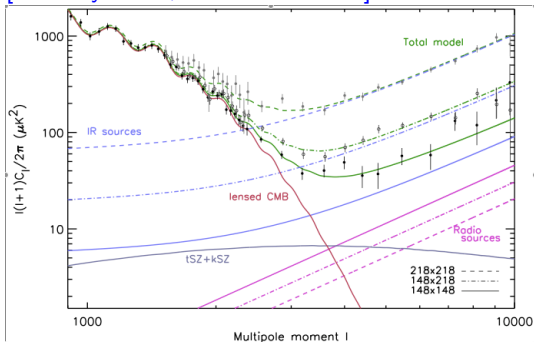
- primary anisotropies negligible due to Silk damping: sets scale for thickness of last-scattering surface
- non-linear effects from more recent epochs: secondary sources of perturbations from primordial potential–matter coupling.



[ACT research proposal]

Information measured in the CMB

[Dunkley et al., arXiv:1009.0866]



- **thermal Sunyaev-Zel'dovich**
⇒ spectral distortion of photons to higher energies.
- **Ostriker-Vishniac (cosmological velocity) and kinetic Sunyaev-Zel'dovich (peculiar motion of cluster)**
⇒ velocity–density coupling.
- **gravitational lensing** ⇒ smearing of acoustic peaks
+ generation of small scale power at temperature

Questions call for more experiments

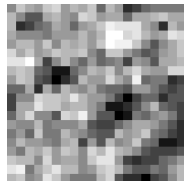
To answer questions such as:

- Did cosmic structure form solely via gravitational instability?
- What is the mass of the neutrino and how does it affect the structure formation?
- How does galaxy distribution relate to the mass distribution?
- What is the nature of the dark matter and the dark energy?
- Are we missing any important ingredient from the cosmological model?

need experiments that can probe the corresponding scales.

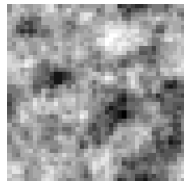
WMAP: 15'/pixel

[<http://lambda.gsfc.nasa.gov/>]



Planck space probe: 7.2'/pixel

[<http://www.rssd.esa.int/index.php?project=Planck>]

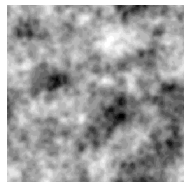


Atacama Cosmology Telescope: 1.4'/pixel

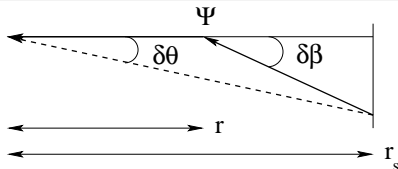
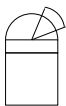
[<http://www.physics.princeton.edu/act/about.html>]

South Pole Telescope: 1.0'/pixel

[<http://pole.uchicago.edu/>]



The lens equation



For a line element

$$ds^2 = a^2(\eta) \left[-(1 - 2\Psi)d\eta^2 + (1 + 2\Psi) \left[dr^2 + f_k^2(r) (d\theta_x^2 + d\theta_y^2) \right] \right]$$

find the lens equation by setting

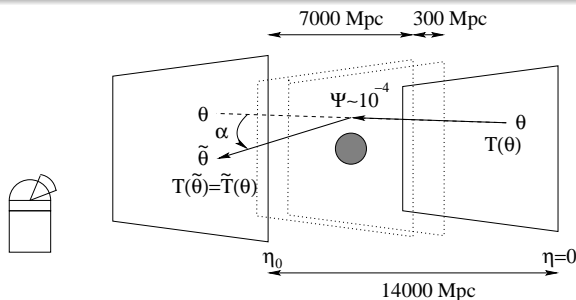
$$f_k(r_s) \delta\theta = f_k(r_s - r) \delta\beta, \text{ where } \delta\beta = -2\delta r [1/f_k(r)] \nabla_{\theta} \Psi(r\theta, r)$$

Adding up the deflections between observer and source

$$\alpha(\theta) = -2 \int_0^{r_s} dr \frac{f_k(r_s - r)}{f_k(r_s) f_k(r)} \nabla_{\theta} \Psi(r\theta, r), \quad f_k(r) = \begin{cases} \sin[r] & : k = +1, \\ r & : k = 0, \\ \sinh[r] & : k = -1. \end{cases}$$

[Challinor & Lewis, PhysRep 429, 1 (2006)]

Lensing of the CMB: probe of large-scale structure



The lensing potential ψ deflects the photons by $\alpha(\theta) = \nabla\psi$
 \Leftrightarrow remap of the temperature anisotropy according to

$$\tilde{T}(\theta) = T(\theta + \nabla\psi) = T(\theta) + \nabla\psi \cdot \nabla T(\theta) + O[(\nabla\psi)^2].$$

14000/300 \sim 50 deflections

Assume random walk:

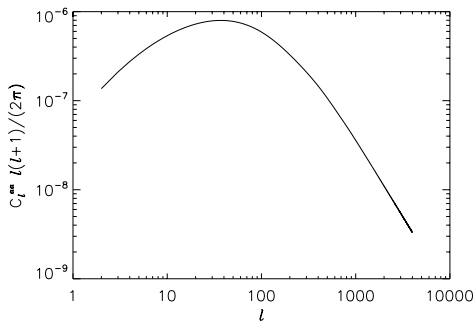
$$\alpha(\theta) \simeq \sqrt{50} \times 10^{-4} \simeq 7 \times 10^{-4} = 2.4 \text{ arcmin}$$

Deflections are coherent over $\theta_{\text{corr}} \simeq 300/7000 \simeq 2 \text{ deg}$

Power spectrum of the lensing deflection

Deflection power spectrum
peaks at $\ell \sim 60 \Rightarrow \theta \sim 3$ deg

$$C_{\ell}^{\alpha\alpha} = [\ell(\ell + 1)]^2 C_{\ell}^{\psi\psi}$$



$$C_{\ell}^{\psi\psi} \approx \frac{8\pi^2}{\ell^3} \int_0^{\eta_0} d\eta \eta \left[\frac{f_k(\eta_0 - \eta)}{f_k(\eta) f_k(\eta_0)} \right]^2 P_{\Psi} \left(\frac{\ell}{f_k(\eta)}, \eta \right)$$

and

$$P_{\Psi}(k, \eta) = \frac{9}{4} \Omega_m^2(\eta) H^4(\eta) \frac{P_{\delta}(k, \eta)}{k^4}$$

Relevance of lensing reconstruction for cosmology

$\hat{\alpha} \times \hat{\alpha}$:

- Subtract lensing B-modes [Seljak & Zaldarriaga, PRL 78, 2054 (1997)]
- Probe density power spectrum [Seljak & Zaldarriaga, PRL 82, 2636 (1999)]
- Constrain sum of neutrino masses and dark energy equation of state [de Putter et al., PRD 79, 065033 (2009); Sherwin et al., PRL 107, 021302 (2011)]

$\hat{\alpha} \times$ Galaxies:

- Estimate galaxy bias [Suginohra et al., ApJ 495, 511 (1998)]

$\hat{\alpha} \times$ Weak lensing:

- Constrain growth of structure [Bean & Tangmatitham, PRD 81 083534 (2010)]
- Constrain distance ratios and curvature [Das & Spergel, PRD 79, 043509 (2009)]

Detections of the lensing signal

- Cross-correlation of WMAP with large-scale structure data:
 $Corr = 1.15 \pm 0.34$ at 3.4σ
[Smith et al., PRD 76, 043510 (2007)]
 $Corr = 1.06 \pm 0.42$ at 2.5σ
[Hirata et al., PRD 78, 043520 (2008)]
- Direct detection from WMAP data:
 $A_L = \{0.96 \pm 0.60, 1.06 \pm 0.69, 0.97 \pm 0.47\}$ at 2.5σ
[Smidt et al., ApJL. 728, 1 (2011)]
- Direct detection from ACT data: $A_L = 1.16 \pm 0.29$ at 4σ
[Das et al., PRL 104, 021301 (2011)]

Estimator of the lensing potential

Lensing effects are apparent in the CMB power spectrum

$$\langle \tilde{T}(\ell') \tilde{T}(\ell' - \ell) \rangle = \delta(\ell) C_{\ell'}^{TT} + \ell \cdot \left[\ell' C_{\ell'}^{TT} + (\ell - \ell') C_{|\ell - \ell'|}^{TT} \right] \psi(\ell)$$

The optimal estimator is the convolution of $\tilde{T}(\ell)$ by $Q^\psi(\ell, \ell')$

$$\hat{\psi}(\ell) = \int d^2\ell' \tilde{T}(\ell') \tilde{T}(\ell - \ell') Q^\psi(\ell, \ell')$$

where

$$Q^\psi(\ell, \ell') = \mathcal{N}_\ell \frac{1}{2} \frac{\ell \cdot \ell' C_{\ell'} + \ell \cdot (\ell - \ell') C_{|\ell - \ell'|}}{\tilde{C}_{\ell'} \tilde{C}_{|\ell - \ell'|}}$$

$$\mathcal{N}_\ell = \left[\int \frac{d^2\ell'}{(2\pi)^2} \frac{1}{2} \frac{[\ell \cdot \ell' C_{\ell'} + \ell \cdot (\ell - \ell') C_{|\ell - \ell'|}]^2}{\tilde{C}_{\ell'} \tilde{C}_{|\ell - \ell'|}} \right]^{-1} \equiv C_\ell^{\hat{\psi}(\ell)}.$$

Do we need a new estimator?

Optimal reconstruction in harmonic space [Hu, ApJ 557, 79 (2001); Okamoto & Hu, PRD 67, 083002 (2003)] implicitly assumes full-sky coverage without:

- galactic cuts,
- bad pixels due to excision of point sources,
- nonuniform weighting for uneven sky coverage.

We consider a slightly less optimal estimator which:

- has compact support,
- acts in real space.

[Carvalho & Moodley, PRD 81, 123010 (2010)]

[Bucher, Carvalho, Moodley & Remazzeilles, arXiv:1004.3285]

A new estimator for the lensing potential

An estimator for the convergence $\kappa_0 = -\frac{1}{2}\nabla^2\psi$ in real space

$$\begin{aligned}\hat{\kappa}_0(\boldsymbol{\theta}) &= \int \frac{d^2\boldsymbol{\ell}}{(2\pi)^2} \exp[i\boldsymbol{\ell} \cdot \boldsymbol{\theta}] \int d^2\boldsymbol{\ell}' \tilde{T}(\boldsymbol{\ell}') \tilde{T}(\boldsymbol{\ell} - \boldsymbol{\ell}') Q(\boldsymbol{\ell}, \boldsymbol{\ell}') \\ &\equiv \int d^2\boldsymbol{\theta}' \tilde{T}(\boldsymbol{\theta}') \int d^2\boldsymbol{\theta}'' \tilde{T}(\boldsymbol{\theta}'') Q(\boldsymbol{\theta}, \boldsymbol{\theta}', \boldsymbol{\theta}''),\end{aligned}$$

where we define the weight function in real space by

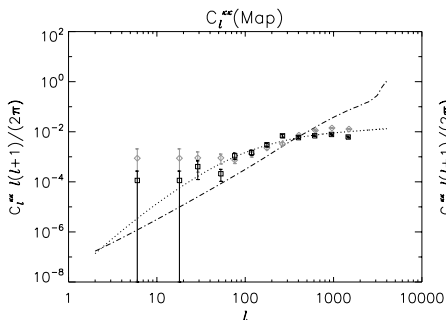
$$\begin{aligned}Q(\boldsymbol{\theta}, \boldsymbol{\theta}', \boldsymbol{\theta}'') &= \int \frac{d^2\boldsymbol{\ell}}{(2\pi)^2} \exp[i\boldsymbol{\ell} \cdot \boldsymbol{\theta}] \\ &\times \int \frac{d^2\boldsymbol{\ell}'}{(2\pi)^2} \exp[-i\boldsymbol{\ell}' \cdot \boldsymbol{\theta}'] \exp[-i(\boldsymbol{\ell} - \boldsymbol{\ell}') \cdot \boldsymbol{\theta}''] Q(\boldsymbol{\ell}, \boldsymbol{\ell}').\end{aligned}$$

and $Q(\boldsymbol{\ell}, \boldsymbol{\ell}') = \ell^2 Q^\psi(\boldsymbol{\ell}, \boldsymbol{\ell}')$.

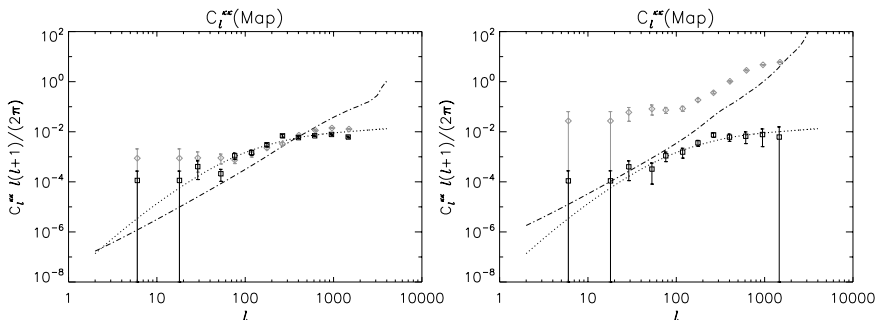
[Carvalho & Moodley, PRD 81, 123010 (2010)]

Estimated $C_\ell^{\kappa\kappa}$ for a Planck simulation

Noiseless



Noisy



- Loss of power at small scales: in theory $\ell_{max} \sim 1/\theta_{fwhm}$ but in practice limited by $1/\theta_{kernel}$ which measures the smallest wavelengths probed by the kernel
- No sensitivity to experimental noise: test with white noise only \Rightarrow noise, independent in each pixel, is averaged out.

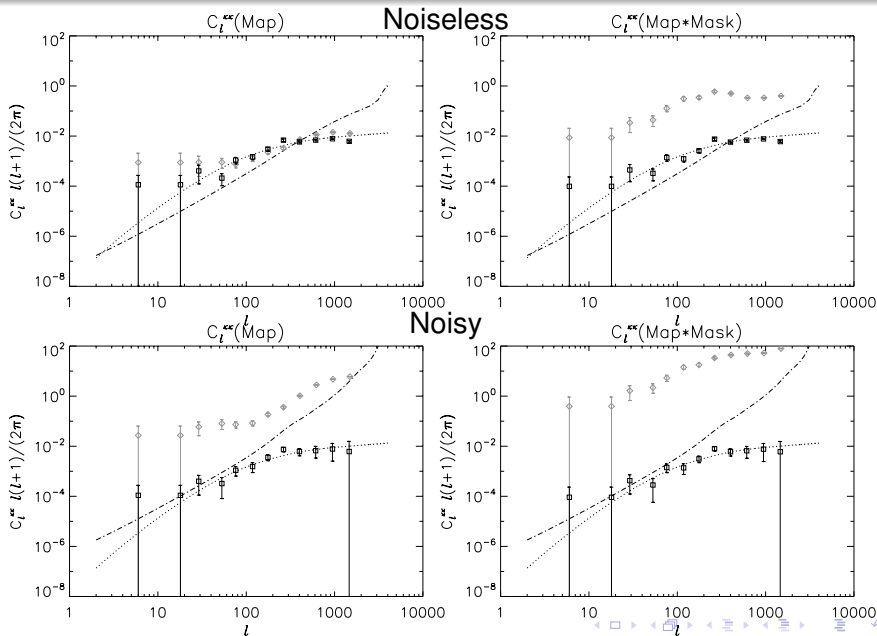
Application to maps with excisions of points

- Foreground signals can generate spurious correlations.
- Weak lensing induces non-Gaussianities in the form of mode couplings \Rightarrow sensitive to foreground signals.

To which extent does masking small areas at the locations of point sources affect the weak lensing convergence extracted from CMB maps? [Carvalho & Tereno, PRD 84, 063001 (2011)]

[Lensed Map \times Beam+Noise] \times





Lensing: probe of dark matter and dark energy

- Convergence of the CMB:

$$\kappa(\hat{n}) = \frac{3}{2} \Omega_m H_0^2 \int_0^{\eta_0} d\eta \frac{\eta}{a(\eta)} \frac{\eta_0 - \eta}{\eta_0} \delta(\eta \hat{n}, \eta)$$

dependent on total mass

⇒ probes statistical properties of large scale structure

- Convergence of distant galaxies:

$$\kappa(\hat{n}) = \frac{3}{2} \Omega_m H_0^2 \int_{\eta_{gal}}^{\eta_0} d\eta \frac{\eta}{a(\eta)} g(\eta) \delta(\eta \hat{n}, \eta)$$

$$g(\eta) = \int_{\eta_{gal}}^{\eta} d\eta' (dN/d\eta') \frac{\eta' - \eta}{\eta'}$$

dependent on (total mass + source redshift distribution)

⇒ probes dark energy properties.

- affects distant-redshift relation
- suppresses time-dependent growth of structures

Lensing: probe of modifications of gravity

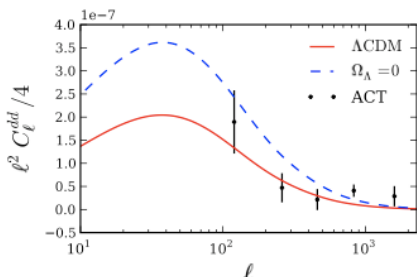
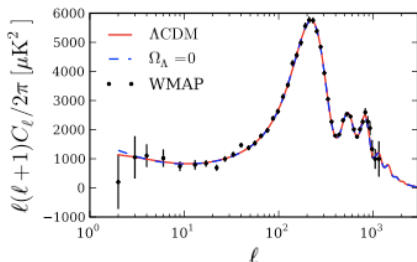
Modified gravity must converge to General Relativity:

- at high z , when expansion is not accelerated
- at small scales, where GR is well tested.

To distinguish between effects of MG and (GR+dark energy):

- impossible with background expansion history only
- need also growth history.

Weak lensing probes the growth
 \Rightarrow can be used to distinguish between models.



[Blake et al., PRL 107, 021301 (2011)]

Parametrization of modified gravity

Parametric Post Friedmann prescription: a model independent description of deviations from GR:

- gravitational split function $\phi - \psi = \zeta\phi + \boxed{g/k} \dot{\phi}$
- mass screening function: $G_{eff} = G\mu^{-2} = G(1 + H \boxed{g/k})$

[Ferreira & Skordis, PRD 81, 1045020 (2010)]

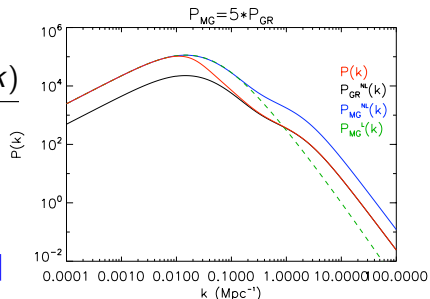
At small scales expect deviations from the linear regime of the matter power spectrum

$$P_{HS}(k) = \frac{P_{MG}^{nl}(k) + c_{nl}\Sigma^2(k)P_{GR}^{nl}(k)}{1 + c_{nl}\Sigma^2(k)}$$

$\Sigma^2(k)$: picks up non-linear scales

c_{nl} : defines scale $P_{HS} \rightarrow P_{GR}$

[Hu & Sawicki, PRD 76, 104043 (2007)]



Constrain models of modified gravity with the lensing of galaxies, the **cosmic shear**

$$\kappa_+ = -\frac{1}{2}(\partial_x^2 + \partial_y^2)\psi, \quad \kappa_\times = -\partial_x\partial_y\psi,$$

and

- Test on COSMOS data: shear of galaxies over $z \in [0, 5]$
[Schrabback et al., A&A 516 (2010)]
- Prepare for the space survey EUCLID: will measure the shape and spectra of galaxies over the entire extragalactic sky in the visible and NIR, out to $z = 2$
[<http://www.ias.u-psud.fr/imEuclid>]
- Combine with lensing of the CMB

Summary

- The CMB reaches us from the surface of last scattering
⇒ contains information about the early Universe and its subsequent evolution.
- Lensing of the CMB imprints in the primordial fluctuations the cumulative deflection by the large-scale structure along the line of sight
⇒ on its own or combined with other effects contributes to clearing out the picture of the Universe.
- A novel approach to extract the lensing field from CMB maps in real space appears to be a promising tool to deal with noise, bias, point source excisions.

Current studies:

- constrain modified gravity models with weak lensing, both cosmic and galactic