

ON THE FLARING GAMMA-RAY ACTIVITY OF QUASAR 3C 279



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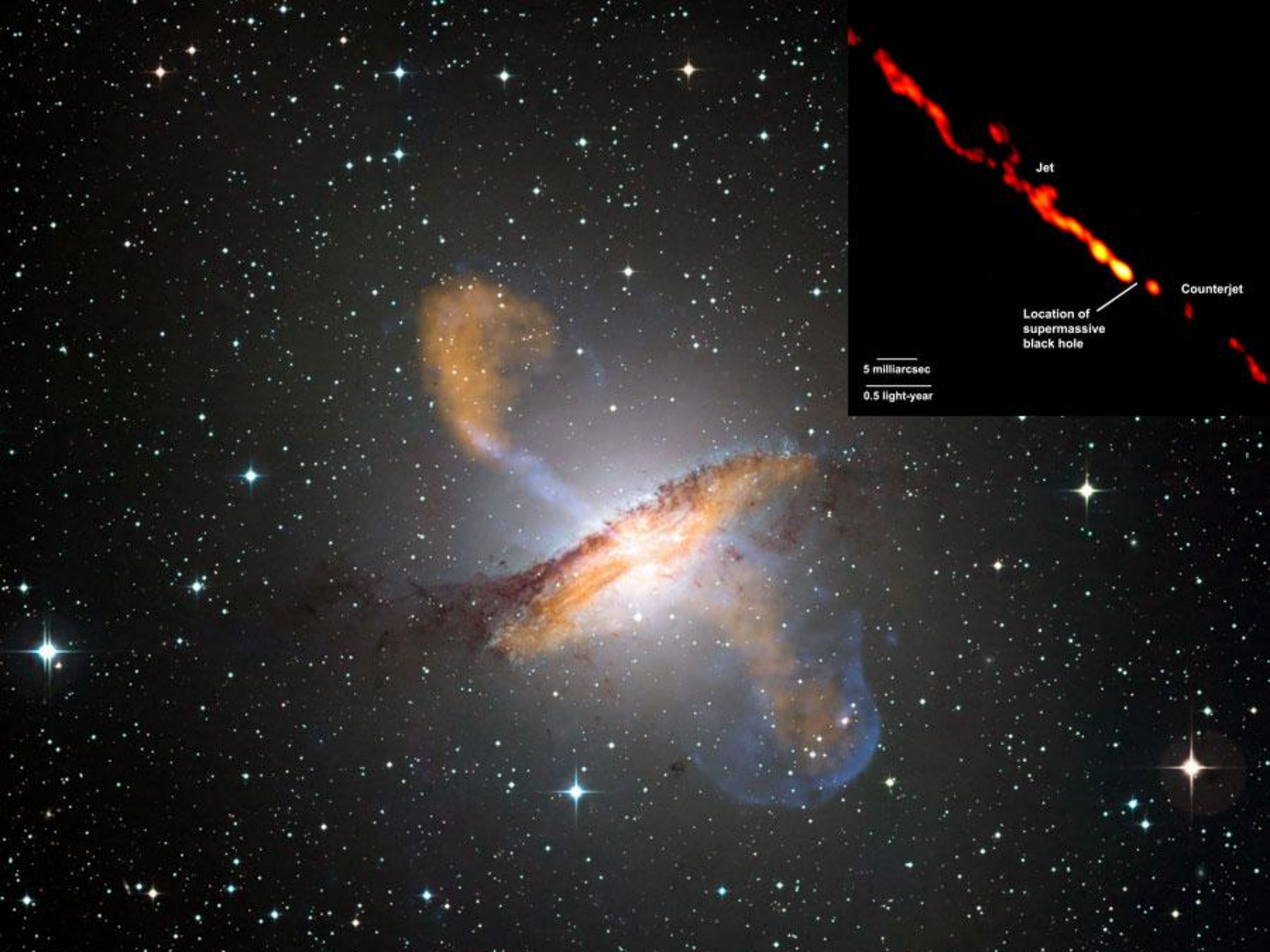
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University of Athens*

TALK OUTLINE

- Active Galactic Nuclei → Blazars
- Models for Multiwavelength emission
- 3C 279: The 'extreme' flare of June 2015

In collaboration with

- Maria Petropoulou – *Purdue University*
- Ioulia Florou – *NKUA*



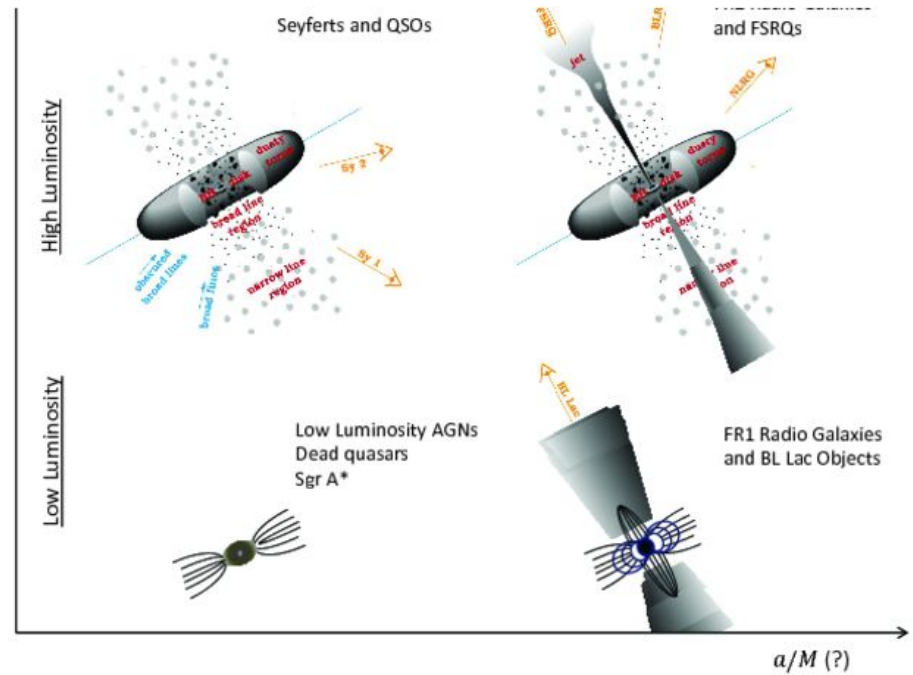
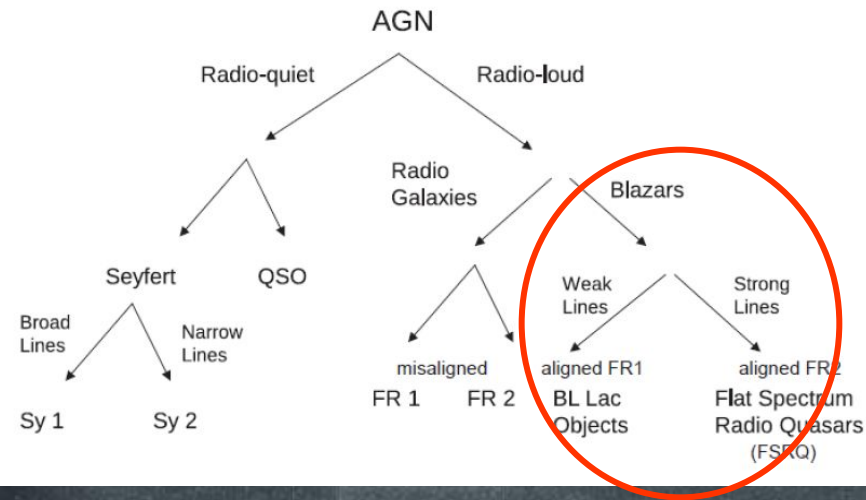
Jet

Counterjet

Location of
supermassive
black hole

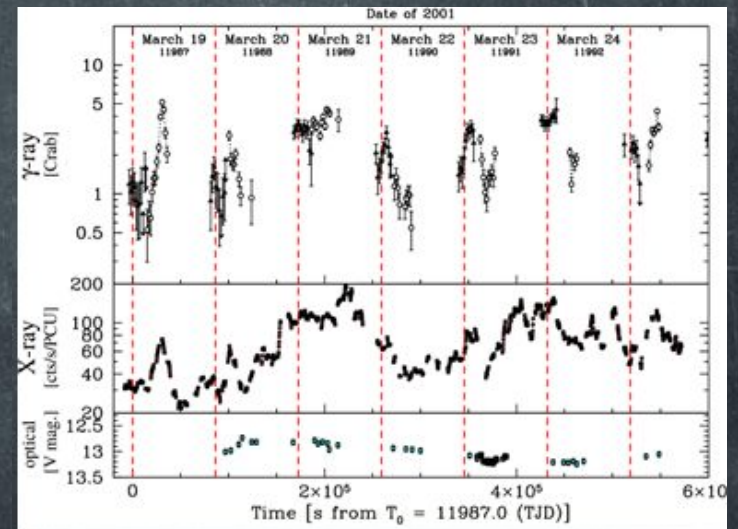
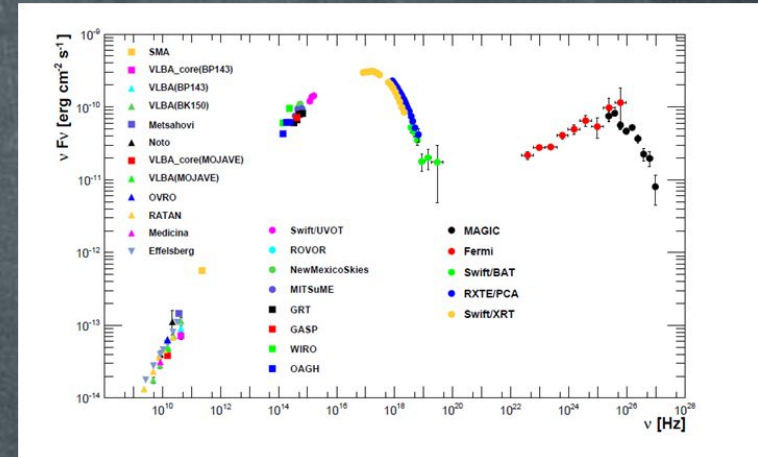
5 milliarcsec
0.5 light-year

AGN TAXONOMY



BLAZARS

- Few (~5% of all AGN)
- Compact , flat spectrum radio sources
- Broad (radio-gamma) non-thermal continuum : 'Double hump' spectrum
- Variable at all energies: short – large amplitude variability + correlations
- Superluminal motion



SPECTRAL FORMATION

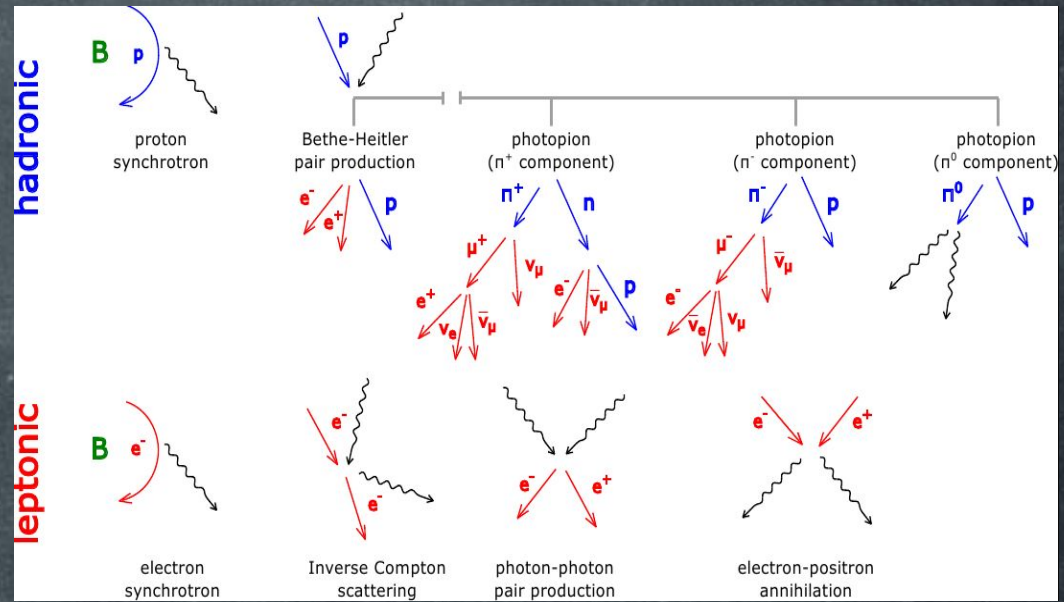
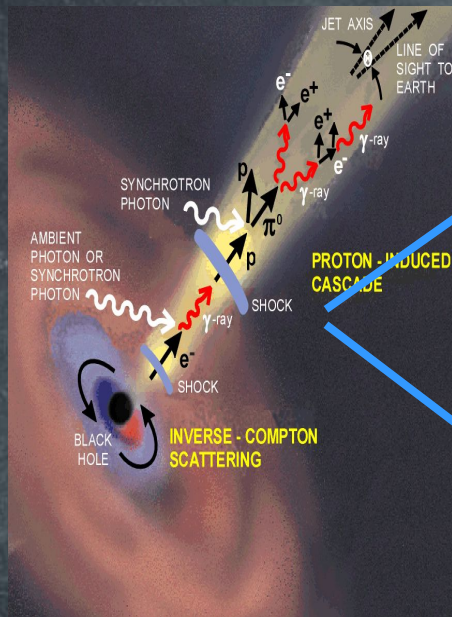
Three observational facts

- MW Spectrum → Particle acceleration to high energies + non-thermal radiation mechanisms
- Fast + correlated variability → emission from a localized region
- Superluminal motion → emission from inside the jet → relativistic beaming

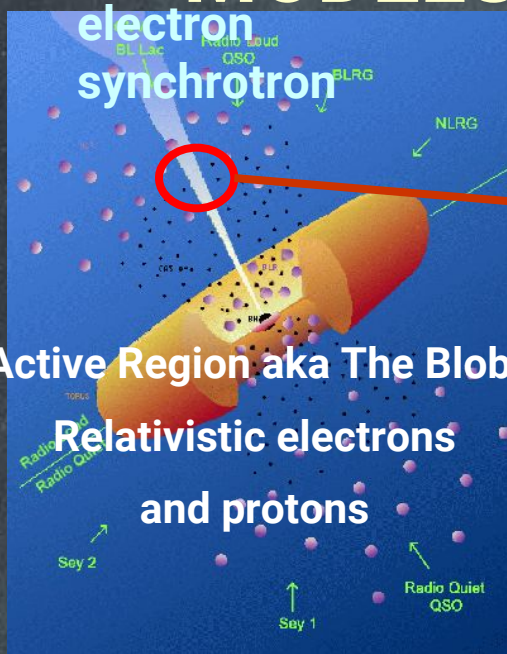
Radiation transfer problem:
theoretical questions

- (1) Geometry of emitting region
- (2) Photon emission + absorption mechanisms
- (3) Particle energization/acceleration
- (4) Species of radiating particles
- (5) Location of 'active' region

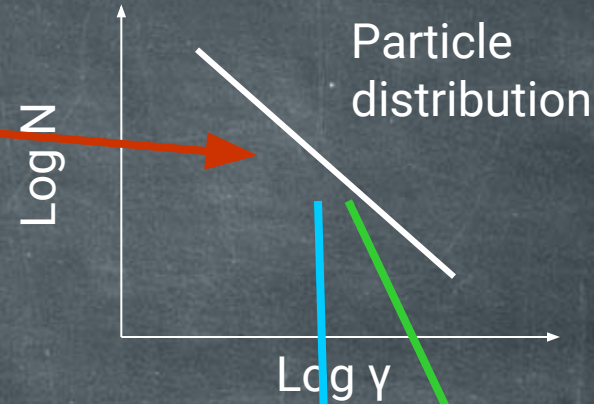
RADIATION MECHANISMS



MODELS FOR BLAZAR EMISSION

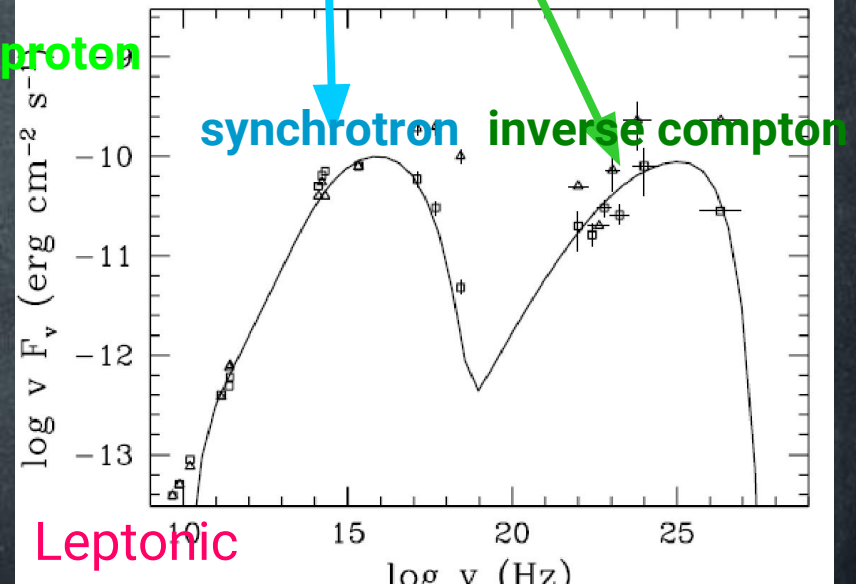
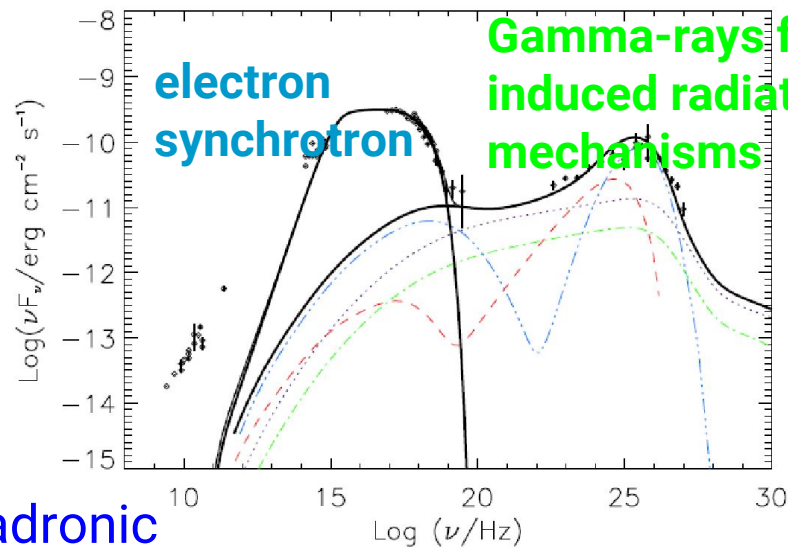


Estimation of B , R_b , δ + specifics of radiating particles

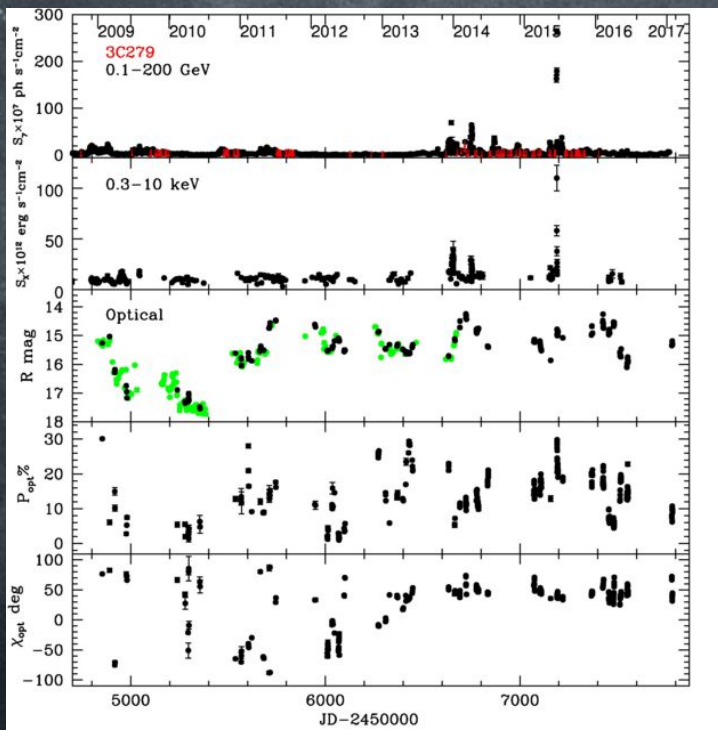
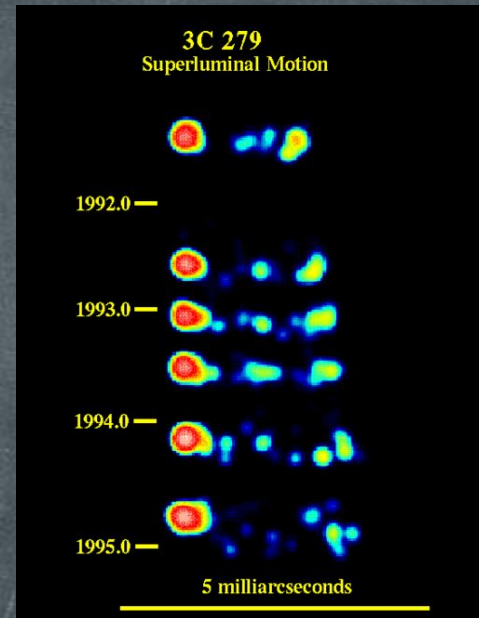
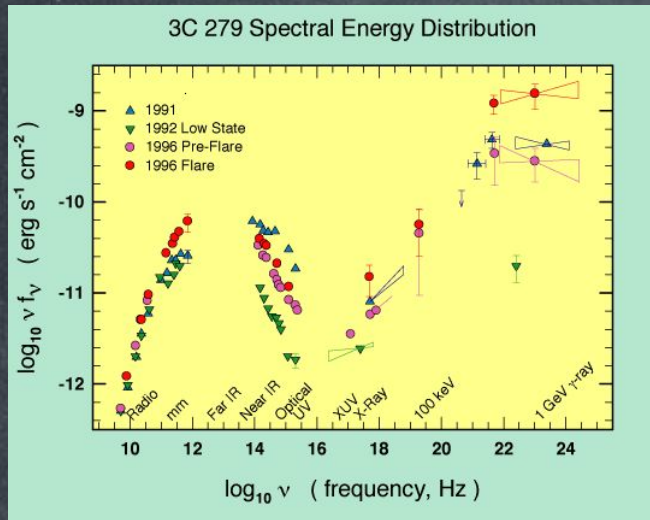


B-field

soft photons



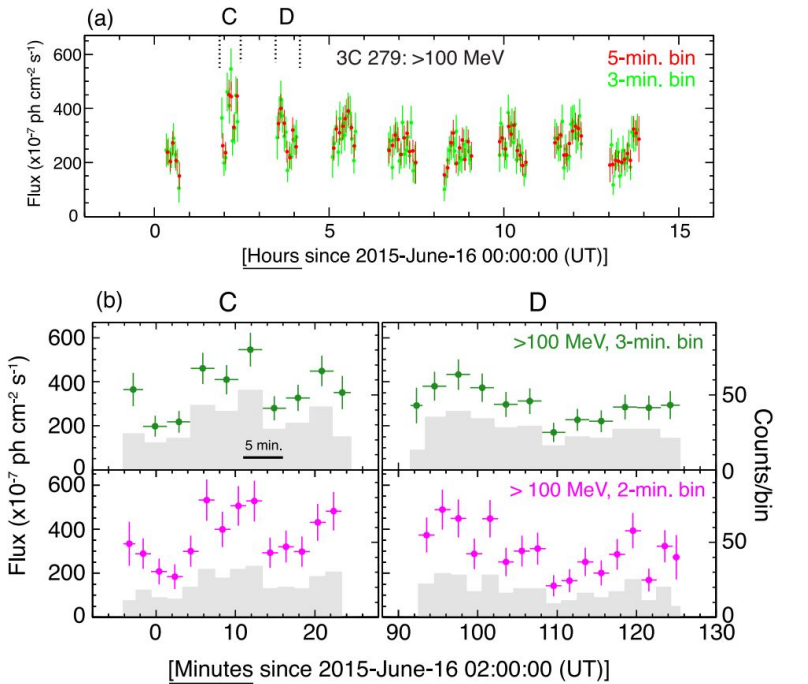
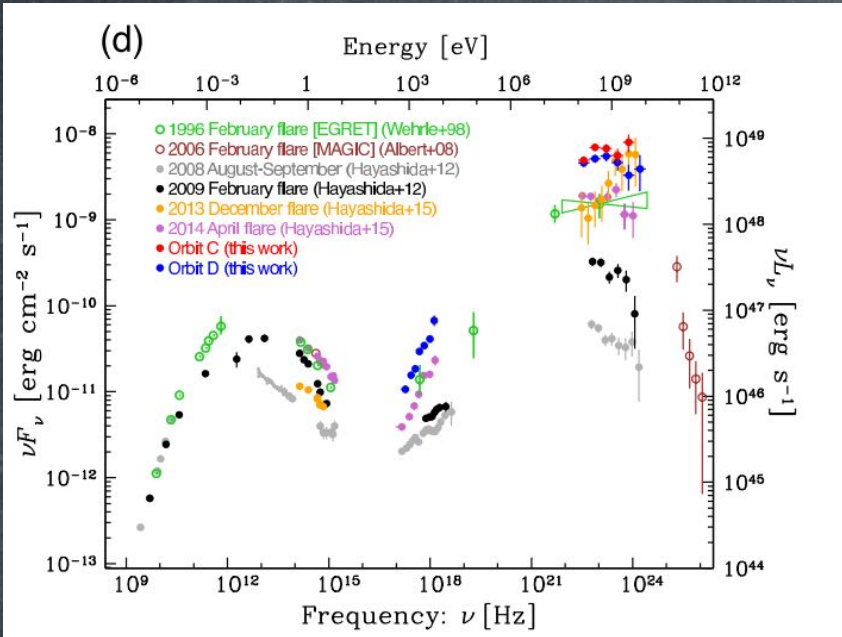
QUASAR 3C 279



Luminosity $\sim 1.e48$ erg/sec

Black hole mass $\sim 5.e8 M_{\text{sun}}$

THE JUNE 2015 GAMMA-RAY FLARE



Giant flare $L_{\gamma} \sim 1.e49$ erg/sec

Subflares as fast as $t_{\text{var}} \sim 5$ min (first ever for Fermi)

A very unique and extreme case for ALL models

THE COMPACTNESS PROBLEM

- Absorption mechanism for gamma-rays on soft photons
- Threshold: if

$$x_\gamma = \frac{h\nu_\gamma}{m_e c^2}$$

then

$$x_{soft} \geq x_{thr} = \frac{2}{x_\gamma}$$

- Optical depth $\tau_{\gamma\gamma} \approx \sigma_{\gamma\gamma} n_s R_b$
- Fast variability \rightarrow small source (or large δ)

$$R_b \sim \delta t_{var} c$$

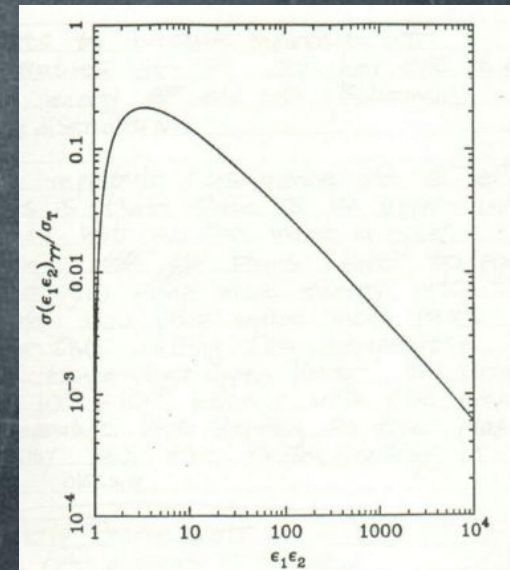
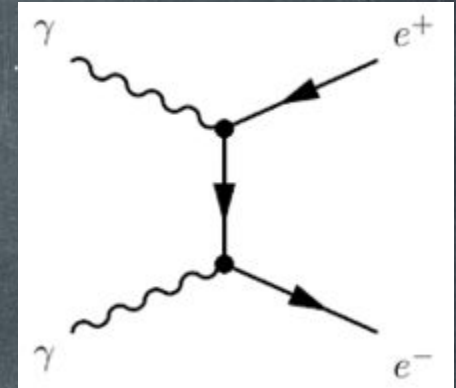
- Target photon density

$$n_s = \frac{L}{4\pi R_b^2 c h \nu_{soft}}$$

therefore

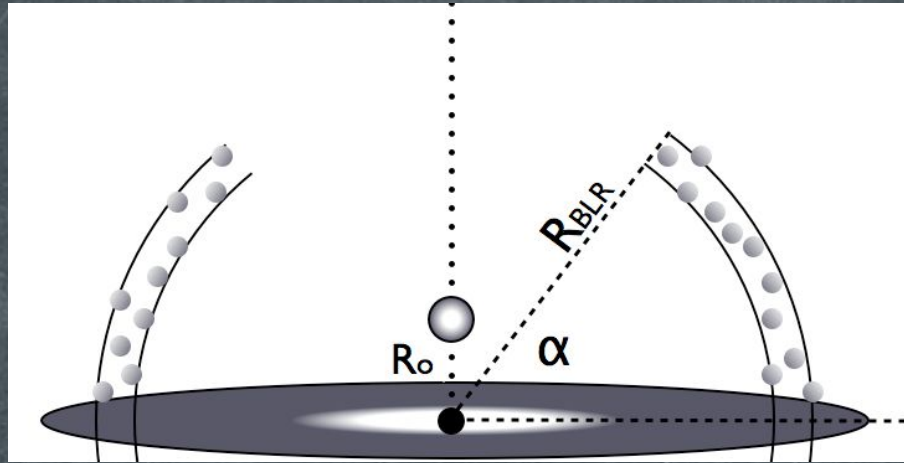
$$\tau_{\gamma\gamma} \sim \frac{L}{R_b}$$

(compactness)



**If compactness high, $\tau_{\gamma\gamma} \gg 1$
 \rightarrow efficient gamma ray
 absorption**

EXTERNAL $\gamma\gamma$ ABSORPTION

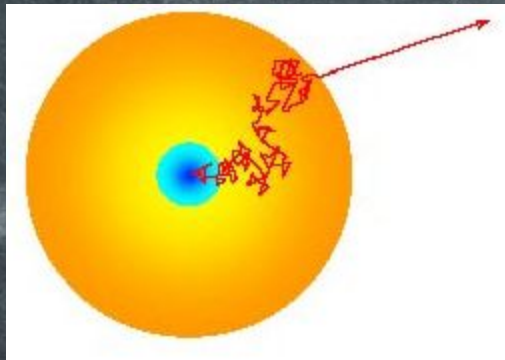


- FSRQs have strong Broad Lines \rightarrow high densities of soft (target) photons
- If $z_0 < R_{BLR} \rightarrow \tau_{\gamma\gamma}$ can be high (“external” absorption)
- Site of radiation zone is crucial for FSRQs

RADIATIVE TRANSFER: AN ANALOGY

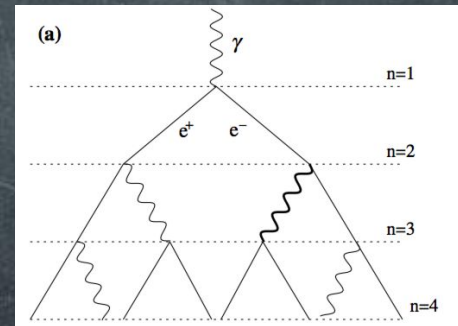
STELLAR ASTROPHYSICS

- Stellar core emits hard photons ϵ_γ (nuclear lines)
- Energy is degraded as photons diffuse on stellar envelope
- Surface emits $\epsilon_\nu \sim kT \ll \epsilon_\gamma$



HIGH ENERGY ASTROPHYSICS

- Blob emits gamma-rays
- If absorbed \rightarrow creation of secondary ee pairs
- Pairs emit more gamma rays
 \rightarrow photon energy is degraded (initial energy is shared by many)
 \rightarrow **electromagnetic cascade**



Luminosity is conserved **but**
photon energy is downgraded

Protons:

$$\frac{\partial n_p}{\partial t} + L_p^{\text{BH}} + L_p^{\text{photonion}} + L_p^{\text{psyn}} + \frac{n_p}{t_{p,\text{esc}}} = Q_p^{\text{inj}} + Q_p^{\text{photonion}}$$

Electrons:

$$\frac{\partial n_e}{\partial t} + L_e^{\text{syn}} + L_e^{\text{ics}} + L_e^{\text{ann}} + L_e^{\text{tpp}} + \frac{n_e}{t_{e,\text{esc}}} = Q_e^{\text{ext}} - Q_e^{\text{BH}} - Q_e^{\gamma\gamma} + Q_e^{\text{photonion}} + Q_e^{\text{tpp}}$$

Photons:

$$\frac{\partial n_\gamma}{\partial t} + \frac{n_\gamma}{t_{\gamma,\text{esc}}} - L_\gamma^{\gamma\gamma} + L_\gamma^{\text{ssa}} = Q_\gamma^{\text{syn}} + Q_\gamma^{\text{psyn}} + Q_\gamma^{\text{ics}} + Q_\gamma^{\text{ann}} + Q_\gamma^{\text{photonion}}$$

Neutrinos:

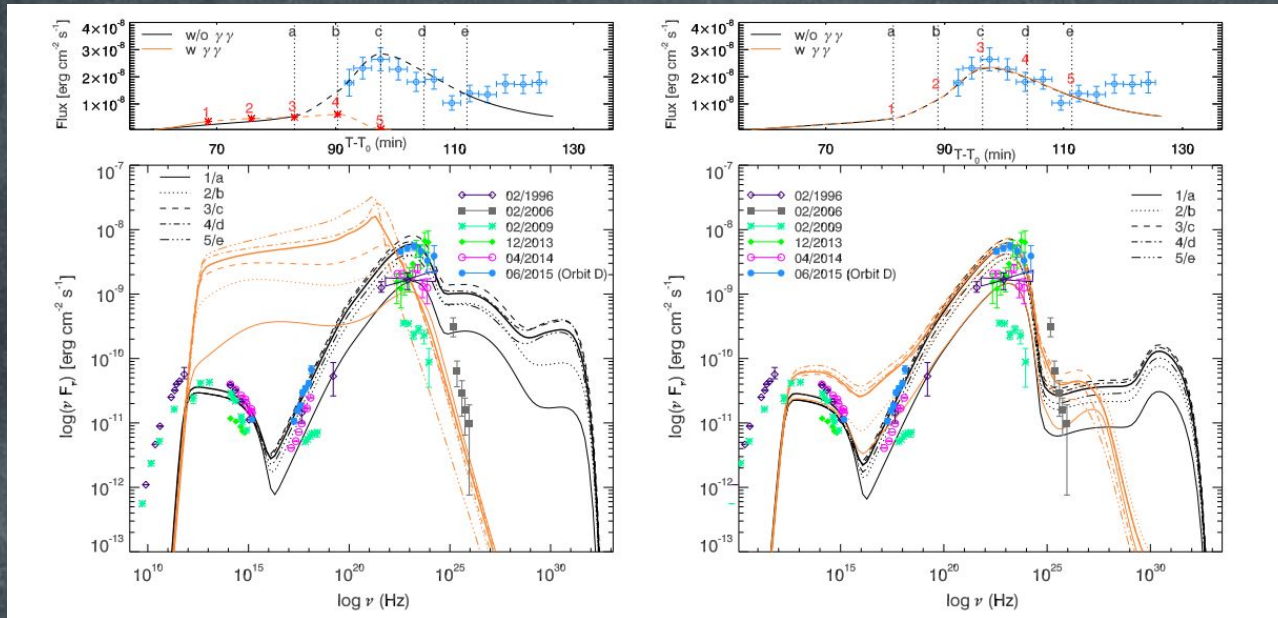
$$\frac{\partial n_\nu}{\partial t} + \frac{n_\nu}{t_{\text{esc}}} = Q_\nu^{\text{photonion}}$$

Neutrons:

$$\frac{\partial n_n}{\partial t} + L_n^{\text{photonion}} + \frac{n_n}{t_{\text{esc}}} = Q_n^{\text{photonion}}$$

- Energy conserving scheme
- Pion, muon & kaon decay is modeled using results of MC code SOPHIA (Muecke et al. 2000)
- Synchrotron cooling of secondaries is also included.

HADRONIC 'FITS' TO THE 3C 279 FLARE



Role of $\gamma\gamma$ absorption crucial

Symbol	Parameter	Case A	Case B
δ	Doppler factor	19.5	50
Γ	Bulk Lorentz factor	9.9	25
B' (kG)	Magnetic field strength	2.2	0.8
r'_b (cm)	Source radius	1.7×10^{14}	4.4×10^{14}
P_j (erg s^{-1})	Absolute total jet power	8.3×10^{47}	10^{48}
u'_p (erg cm^{-3})	Proton energy density	2.6×10^5	7×10^3
u'_B (erg cm^{-3})	Magnetic energy density	2×10^5	2.6×10^4
$E'_{p,\text{max}}$ (eV)	Max. proton energy	5.2×10^{16}	4.8×10^{16}
z_{diss} (cm)	Dissipation distance	1.7×10^{16}	1.1×10^{17}

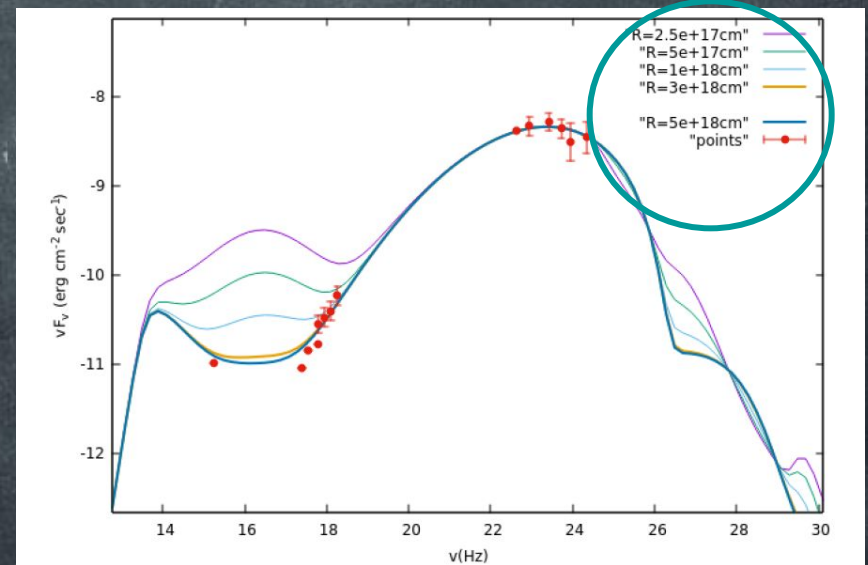
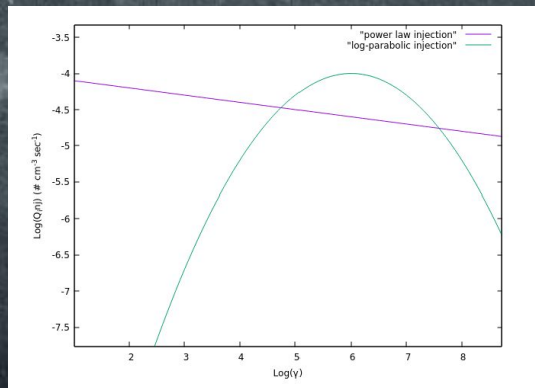
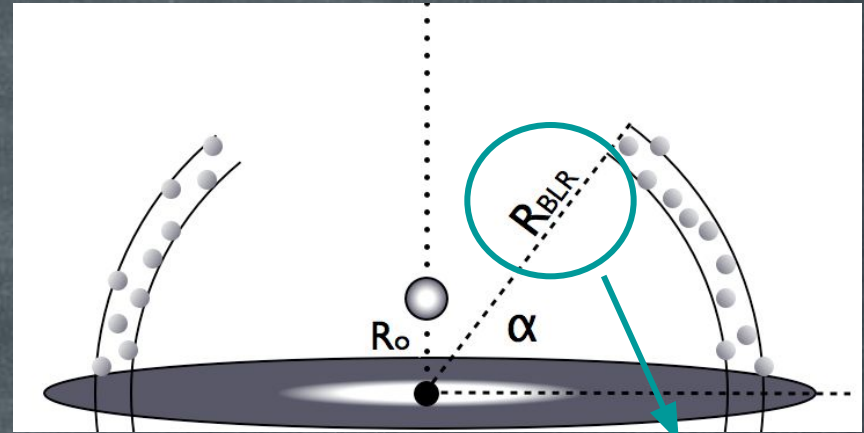
high δ

small source

$10 L_{\text{Edd}}$

ANOTHER EFFORT

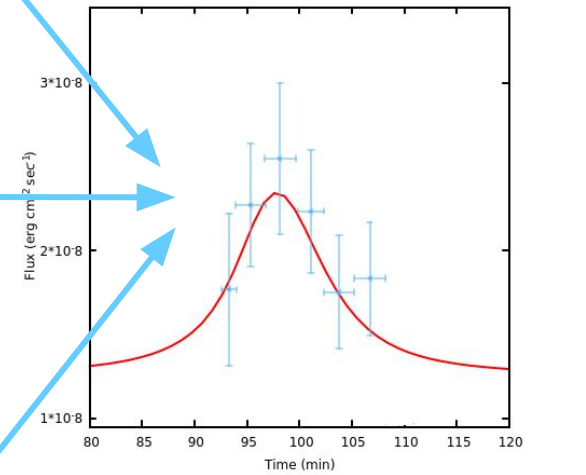
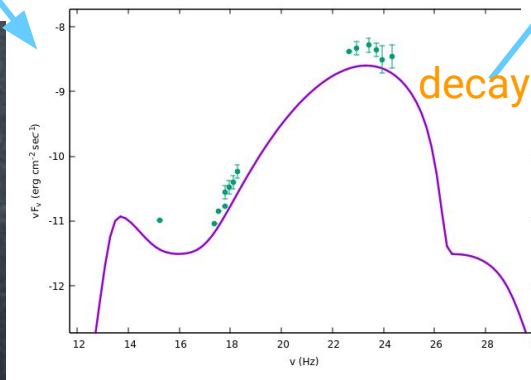
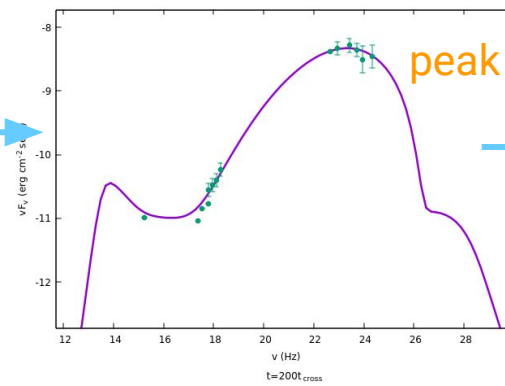
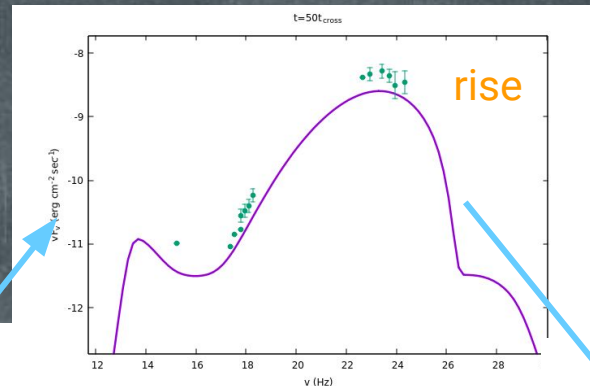
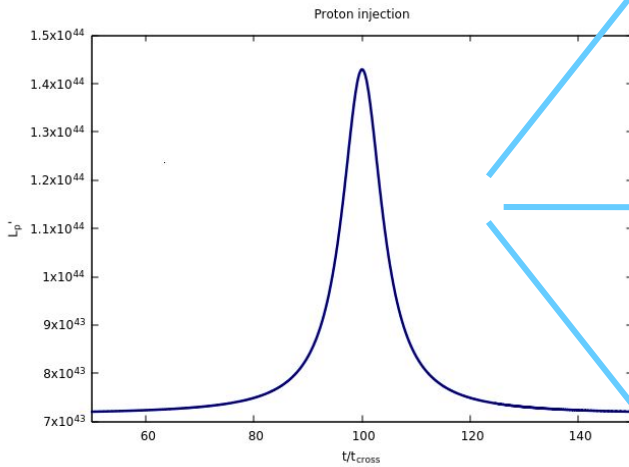
- Change power-law proton injection to log parabolic (more economic)
- Use R_{BLR} as a free (?) parameter
→ minimize external absorption



See Ioulia Florou's poster

MODELING THE 5 min FLARE

$$P_{\text{jet}} \sim 5L_{\text{Edd}}$$



Lorentzian injection
of relativistic protons

Gamma-ray lightcurve

MW photon spectra

See Ioulia Florou's poster

CONCLUSIONS

- The one-zone hadronic model can reproduce (rather) successfully the 'extreme' 5 min gamma-ray flare of 3C 279 of June 2015
- Extreme events require extreme conditions:
 - High δ factors (> 50) – compactness problem!
 - Super Eddington jet power (~ 10) – log-parabolic distributions more economic by a factor of 2.
 - 'Active' region of small size ($\sim 1.e14$ cm) and outside the BLR ($z > 0.1$ pc) – external $\gamma\gamma$ absorption
- Leptonic models?
- Time for a new approach? Maybe more examples needed (it's only a flare, not some steady state)
- Any new model has to treat correctly the radiative transfer for gamma-rays!

BACK UP SLIDES

OVERALL ENERGETICS

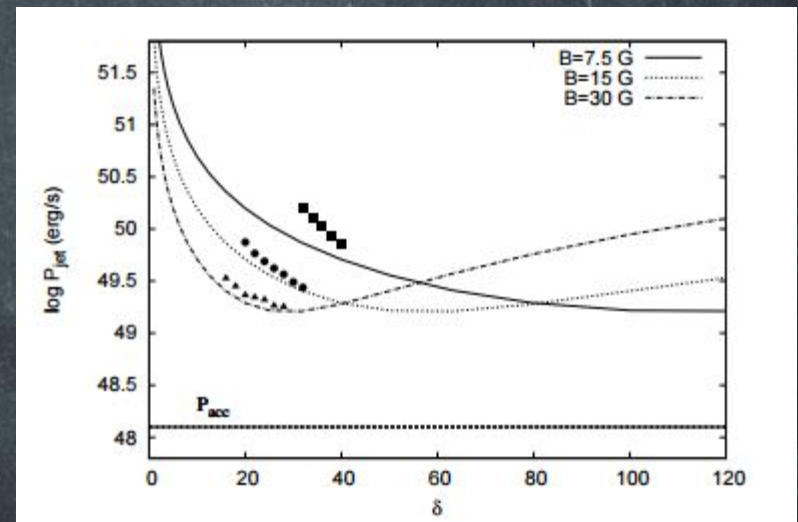
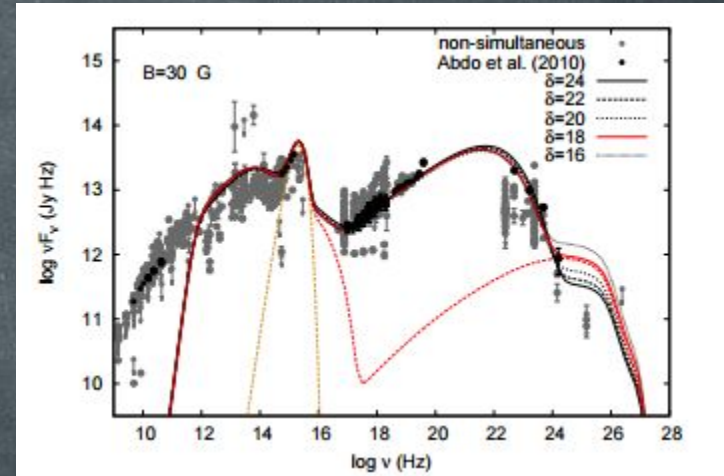
- Simple one-zone synchrotron hadronic fits can be degenerate → different sets of parameters give same fits.
- Minimize the power (similar to equipartition arguments in radio sources with gamma-rays replacing radio and protons replacing electrons) (*Petropoulou & AM 2012*)

$$P_{\text{jet}} \approx \pi R^2 \Gamma^2 c (u_p + u_B)$$

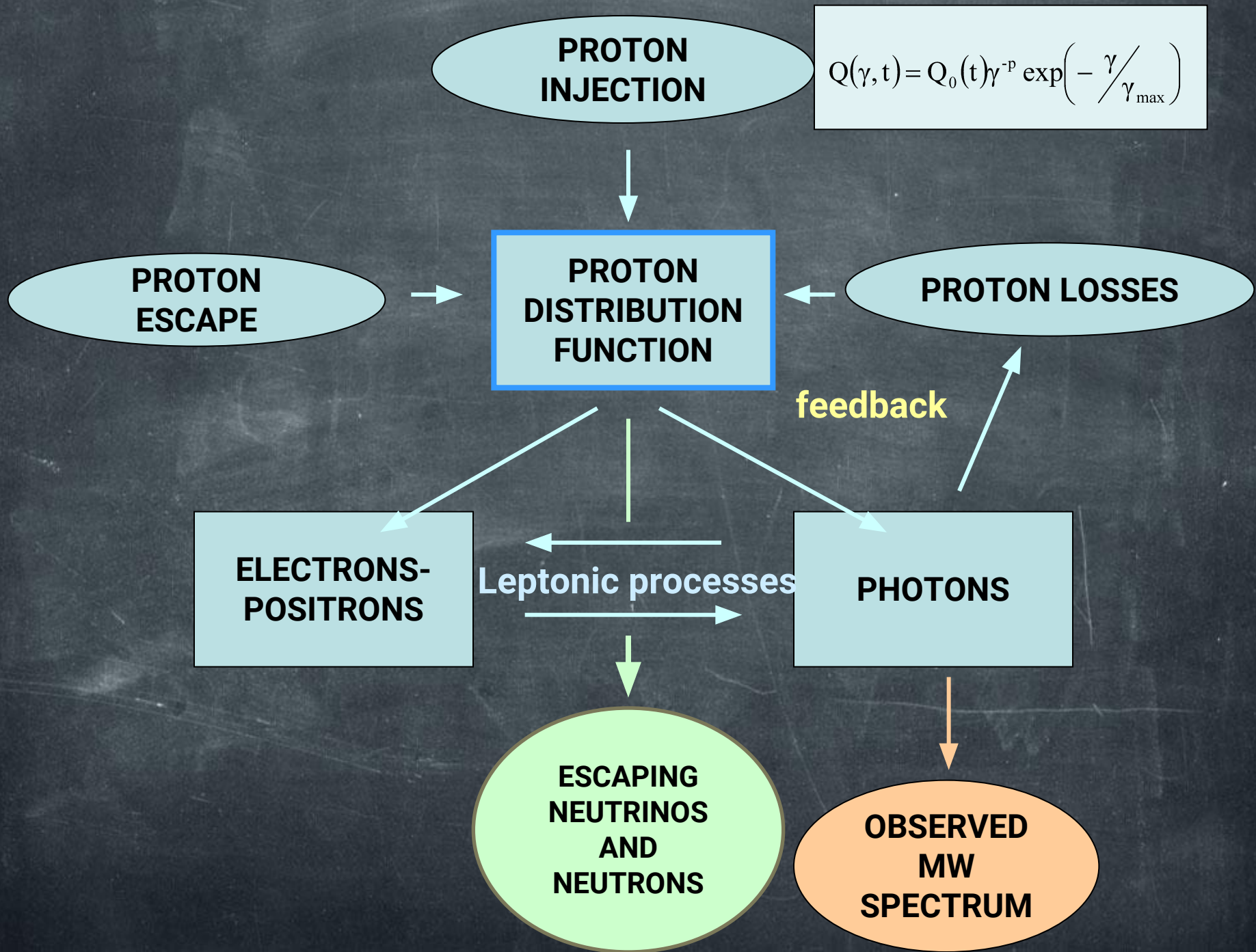
$$P_{\text{jet}} \approx \pi R^2 c \left[A(\delta \cdot B)^{-3/2} + \frac{(\delta \cdot B)^2}{8\pi} \right]$$

$$\frac{dP_{\text{jet}}}{d\delta} = 0 \Rightarrow P_{\text{jet, min}} \text{ for } \delta \cdot B = C$$

3C273



Petropoulou & Dimitrakoudis



A PARADIGM: BOX-MODEL FOR PARTICLE ACCELERATION

$$\frac{\partial N}{\partial t} + \frac{\partial}{\partial E} (r_{\text{acc}} E N - \alpha_s E^2 N) = Q - r_{\text{esc}} N$$

acceleration

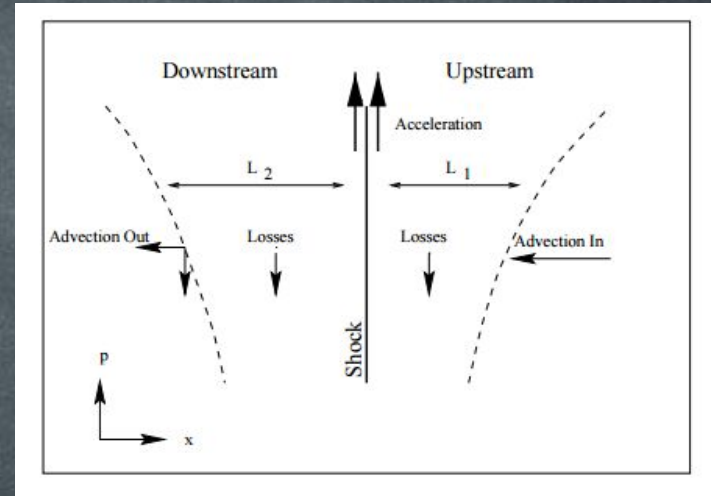
energy loss

escape

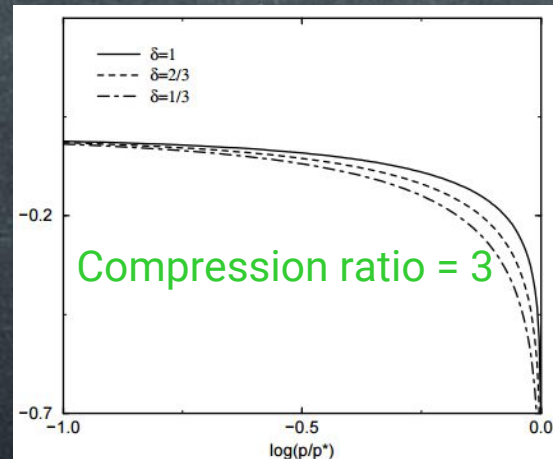
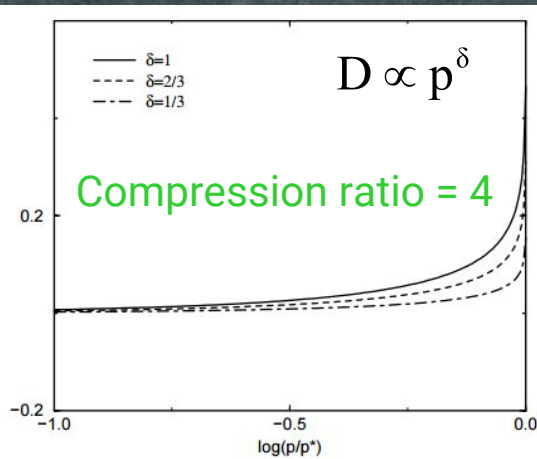
particle injection
at low momenta

$$r_{\text{acc}} = \frac{\bar{n} \cdot (\bar{v}_1 - \bar{v}_2)}{3L}$$

$$r_{\text{esc}} = \frac{\bar{n} \cdot \bar{v}_2}{L}$$



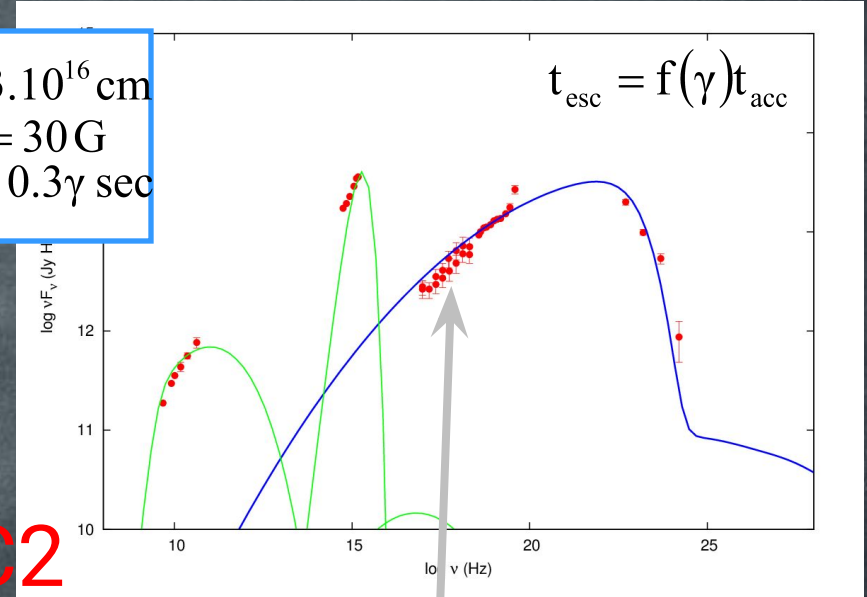
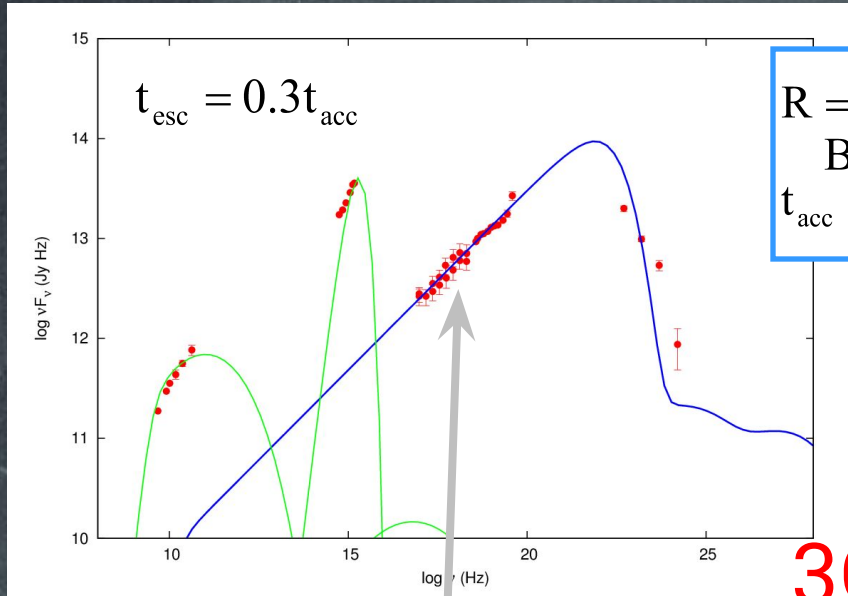
Drury et al. 1999



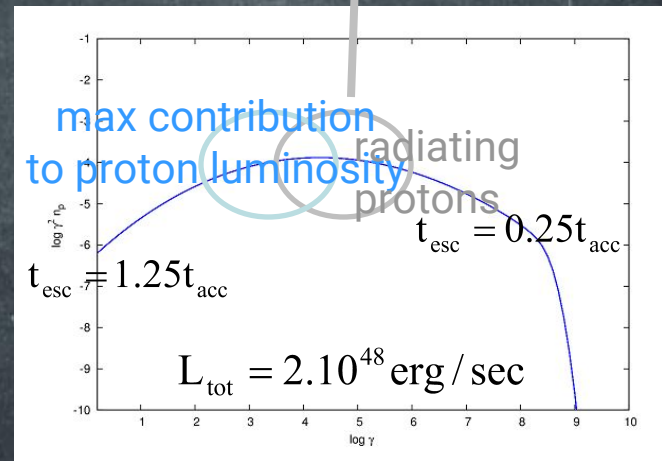
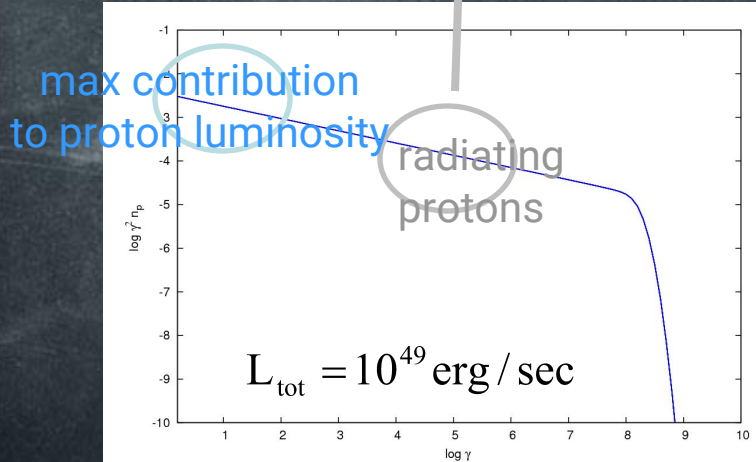
Particle distribution up to a max energy where $T_{\text{acc}} = T_{\text{loss}}$.
Shape of cutoff not *a-priori* assumed (obtained shape non-trivial, e.g. pile-ups)

Standard box model

Modified box model



302
73



$$\frac{\partial N}{\partial t} + \frac{\partial}{\partial E} (r_{\text{acc}} EN) = Q - r_{\text{esc}} N$$

**PROTON
ACCELERATION**

**PROTON
ESCAPE**

**PROTON
DISTRIBUTION
FUNCTION**

PROTON LOSSES

**ELECTRONS-
POSITRONS**

Leptonic processes

PHOTONS

**ESCAPING
NEUTRINOS
AND
NEUTRONS**

**OBSERVED
MW
SPECTRUM**

