



Neutron star magnetospheres with classical and modern methods

5th Hel.A.S Summer School - MHD in Astrophysics

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- ▶ Physical framework
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- ▶ Magnetars with PINNs
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- ▶ Future plans



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Magnetars and Pulsars

1 Introduction

Magnetars and Pulsars

1 Introduction

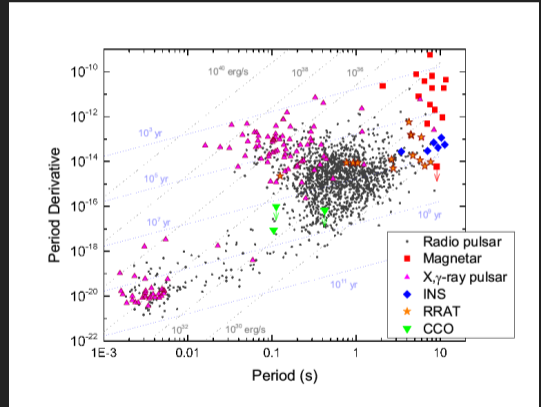


Figure 1: The $P - \dot{P}$ diagram (Harding, 2013)

Magnetars and Pulsars

1 Introduction

Magnetars

- Slow rotation $P \gtrsim 1$ s
- Extreme magnetic field $B \gtrsim 10^{14}$ G
- Young age
- Powered by magnetic field decay

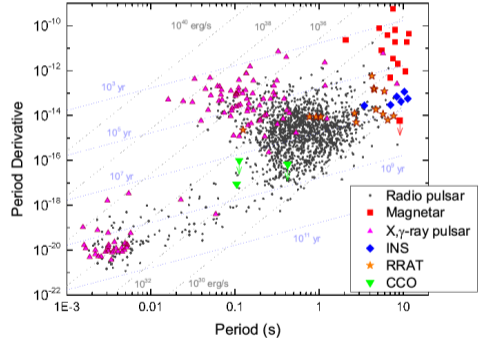


Figure 1: The $P - \dot{P}$ diagram (Harding, 2013)

Magnetars and Pulsars

1 Introduction

Magnetars

- Slow rotation $P \gtrsim 1$ s
- Extreme magnetic field $B \gtrsim 10^{14}$ G
- Young age
- Powered by magnetic field decay

Pulsars

- Fast rotation $P < 1$ s
- Magnetic field $B \sim 10^{12}$ G
- Powered by rotational

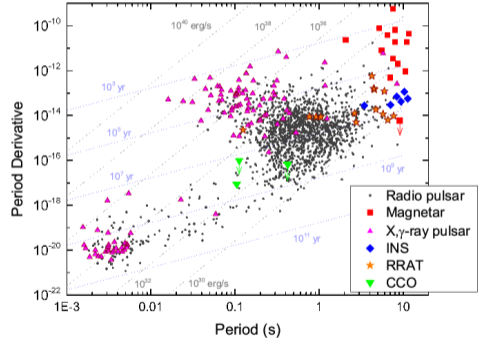


Figure 1: The $P - \dot{P}$ diagram (Harding, 2013)



Observations

1 Introduction



Observations

1 Introduction

Magnetars

- Pulsed X-ray thermal emission
- Transient X-ray activity
 - Short duration bursts
 - Long duration outbursts
 - Highly energetic flares
- Reconfiguration of magnetic field
- Release of energy



Observations

1 Introduction

Magnetars

- Pulsed X-ray thermal emission
- Transient X-ray activity
 - Short duration bursts
 - Long duration outbursts
 - Highly energetic flares
- Reconfiguration of magnetic field
- Release of energy

Pulsars

- Pulsed non-thermal γ -ray emission
- Non-ideal regions
- Acceleration of particles



Magnetosphere

1 Introduction



Magnetosphere

1 Introduction

Definition

The region surrounding the neutron star where plasma flows in its magnetic field



Magnetosphere

1 Introduction

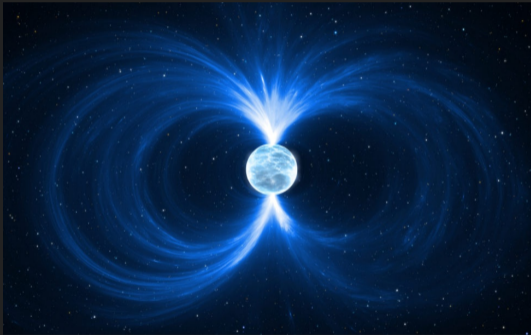


Figure 2: An artist's conception of a NS

Definition

The region surrounding the neutron star where plasma flows in its magnetic field

Magnetosphere

1 Introduction

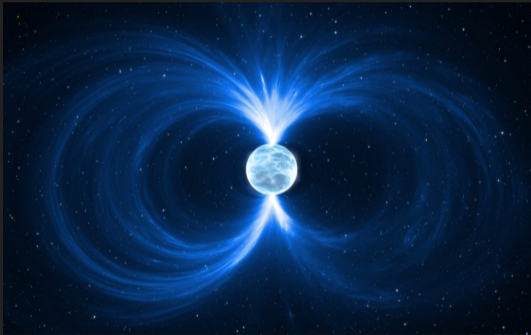


Figure 2: An artist's conception of a NS

Definition

The region surrounding the neutron star where plasma flows in its magnetic field

Characteristics

- Large scale field is roughly dipolar
- Charged particles move along field lines
- Corotation with the star
- Presence of toroidal fields
- Non-thermal emission



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Force-free approximation

2 Physical framework



Force-free approximation

2 Physical framework

Assumptions

- Dynamics is dominated by the electromagnetic field
- Gravity, plasma pressure and particle inertia are negligible
- Quasi-stationary regime

Definitions

$$\alpha = \frac{4\pi}{c} (\mathbf{J} - \rho_e c \boldsymbol{\beta}) \cdot \frac{\mathbf{B}}{B^2}$$
$$\boldsymbol{\beta} = \frac{\boldsymbol{\Omega} \times \mathbf{r}}{c}$$

3D Equations

$$\nabla \times (\mathbf{B} - \beta^2 \mathbf{B}_p) = \alpha \mathbf{B}$$
$$\mathbf{B} \cdot \nabla \alpha = 0$$

Properties

- α is constant along field lines
- Implicitly nonlinear equation
- Equilibrium solutions



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Motivation

3 3D Magnetars

Motivation

3 3D Magnetars

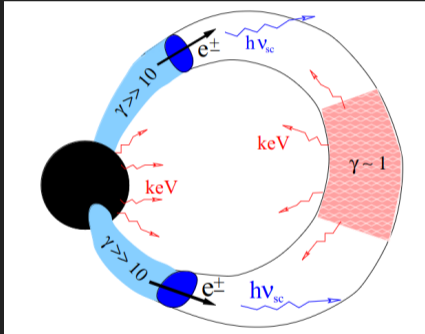


Figure 3: A sketch of a magnetar
(Beloborodov, 2013)

Motivation

3 3D Magnetars

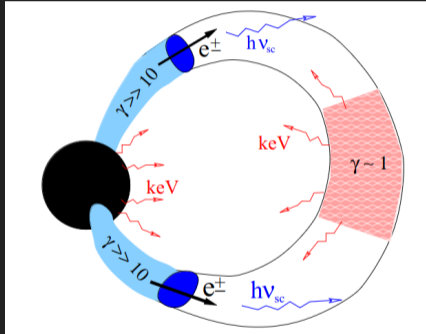


Figure 3: A sketch of a magnetar
(Beloborodov, 2013)

Context

- Internal processes gradually displace field lines
- Energy and helicity flow into the magnetosphere
- Twisted magnetic loops appear in localised regions
- Critical point
- Release of magnetic energy
- Only axisymmetric solutions \rightarrow novelty



Equations

3 3D Magnetars



Equations

3 3D Magnetars

Non-rotating limit

$$\nabla \times \mathbf{B} = \alpha \mathbf{B}$$

$$\mathbf{B} \cdot \nabla \alpha = 0$$

Properties

- $\mathbf{J} \parallel \mathbf{B}$
- Nonlinear system if $\alpha \neq \text{const}$
- Only numerical solutions



The Grad-Rubin method

3 3D Magnetars

Hyperbolic part

$$\mathbf{B}^{(k)} \cdot \nabla \alpha^{(k+1)} = 0$$

- Solves for α with a given \mathbf{B}
- *Characteristics Method*



The Grad-Rubin method

3 3D Magnetars

Hyperbolic part

$$\mathbf{B}^{(k)} \cdot \nabla_{\alpha^{(k+1)}} = 0$$

- Solves for α with a given \mathbf{B}
- *Characteristics Method*

Elliptic part

$$\nabla^2 \mathbf{A}^{(k+1)} = -\alpha^{(k+1)} \mathbf{B}^{(k)}$$

- Solves for \mathbf{A} with a given α
- *Scheduled Relaxation Jacobi method*



Model

3 3D Magnetars

Features

$$\alpha_{S^+}(\theta, \phi) = \alpha_0 \exp \left[\frac{-(\theta - \theta_1)^2 - (\phi - \phi_1)^2}{2\sigma^2} \right]$$

- Two hot spots connected by a twisted loop
- One spot has a Gaussian profile
- Dipole background
- Rest of the magnetosphere is current-free

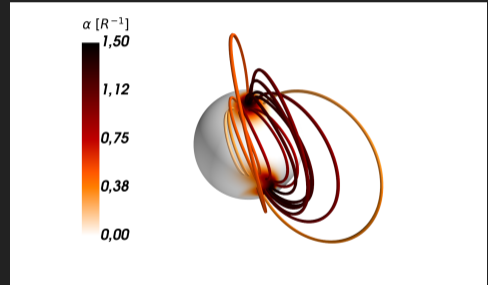


Figure 4: A model with two hotspots



Results

3 3D Magnetars

Results

3 3D Magnetars

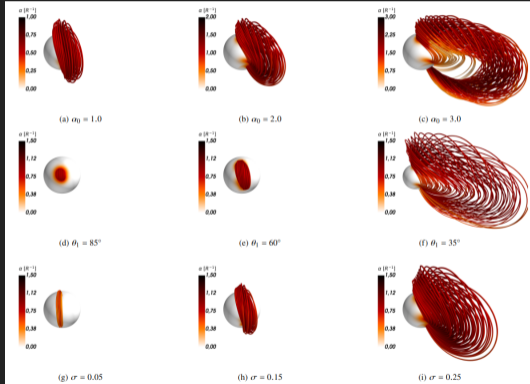


Figure 5: A selection of models for different values of α_0 , θ_1 , σ

Results

3 3D Magnetars

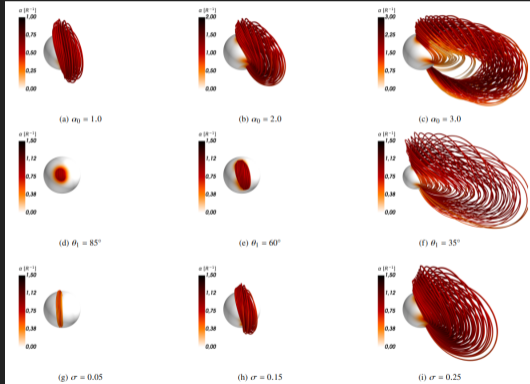


Figure 5: A selection of models for different values of $\alpha_0, \theta_1, \sigma$

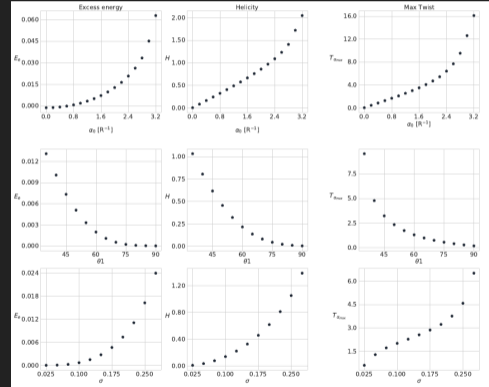


Figure 6: Dependence of energy, helicity and twist on $\alpha_0, \theta_1, \sigma$



Discussion

3 3D Magnetars



Discussion

3 3D Magnetars

Key points

- For certain values of $\alpha_0, \theta_1, \sigma$ the solver does not converge
- Non-existence of solution \rightarrow release of energy
- Excess energy is enough to feed X-ray activity
- Currents heat the surface \rightarrow X-ray pulses

Key points

- For certain values of $\alpha_0, \theta_1, \sigma$ the solver does not converge
- Non-existence of solution \rightarrow release of energy
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- Currents heat the surface \rightarrow X-ray pulses

T_{eff} [keV]
0,24
0,21
0,17
0,14
0,10

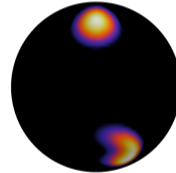


Figure 7: Effective temperature of the hotspots at the surface of the star



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PINNs

4 Magnetars with PINNs

Basics

- Proposed in 1997
- Solve PDEs using neural networks
- Input: a set of points in some domain
- Output: an approximate solution
- Loss: the residuals of the PDE describing the system

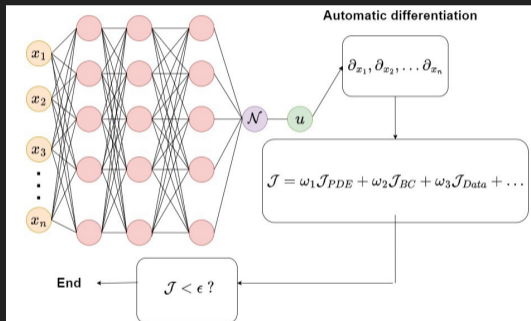


Figure 8: A physics-informed neural network



Motivation

4 Magnetars with PINNs



Motivation

4 Magnetars with PINNs

Relevance

New, promising, exciting, actively developing method for solving PDEs



Motivation

4 Magnetars with PINNs

Relevance

New, promising, exciting, actively developing method for solving PDEs

Advantages

- Meshless
- Scalable
- Generalisable
- Very fast once trained



Motivation

4 Magnetars with PINNs

Relevance

New, promising, exciting, actively developing method for solving PDEs

Advantages

- Meshless
- Scalable
- Generalisable
- Very fast once trained

Disadvantages (for now)

- Precision
- Efficiency
- Rigorous mathematical support
- Relies on trial and error



Motivation

4 Magnetars with PINNs



Motivation

4 Magnetars with PINNs

Context

- PINNs for elliptic problems
- Axisymmetric magnetars
- Evaluate multiple solutions
- Identify possible caveats

Novelty

- Train **once** for a family of solutions
- Treat boundary conditions and source terms as **inputs** to the PINN
- Connect dissimilar domains



Equations

4 Magnetars with PINNs

Equations

4 Magnetars with PINNs

Axisymmetric magnetar magnetospheres

- Compactified polar coordinates $\rightarrow q = \frac{R}{r}, \mu = \cos \theta$
- Axisymmetry $\rightarrow \mathbf{B} = \frac{q}{\sqrt{1-\mu^2}} (\nabla \mathcal{P} \times \hat{\phi} + \mathcal{T} \hat{\phi})$
- Grad-Shafranov equation

$$q^2 \frac{\partial}{\partial q} \left(q^2 \frac{\partial \mathcal{P}}{\partial q} \right) + (1 - \mu^2) q^2 \frac{\partial^2 \mathcal{P}}{\partial \mu^2} + \mathcal{T} \frac{d\mathcal{T}}{d\mathcal{P}} = 0$$

- The solution is completely determined by the surface boundary condition

$$\mathcal{P}(q = 1, \mu) = (1 - \mu^2) \sum_{l=1}^{l_{\max}} \frac{b_l}{l} P'_l(\mu)$$

Results

4 Magnetars with PINNs

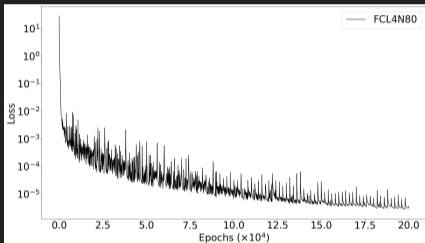


Figure 9: Loss function vs. iterations

	E_{φ}	E_{B_r}	$E_{B_{\theta}}$	E_B
L_1 norm	0.017	0.017	0.025	0.015
L_2 norm	0.019	0.023	0.045	0.023

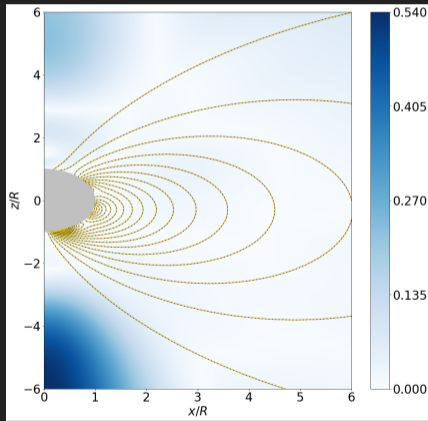


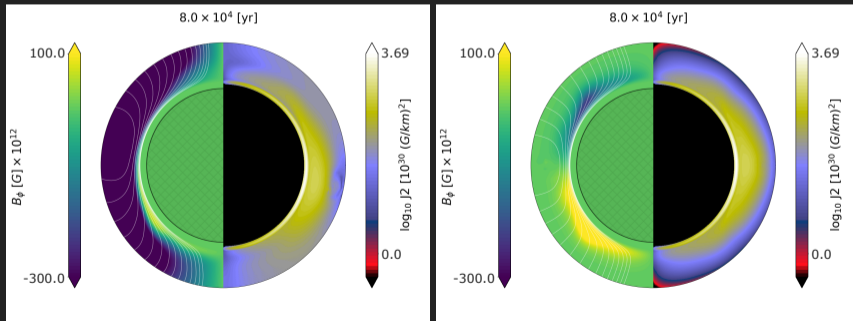
Figure 10: Contours of P

Results

4 Magnetars with PINNs

Application

- Couple magnetothermal evolution with a force-free magnetosphere
- Currents can thread the surface





Discussion

4 Magnetars with PINNs



Discussion

4 Magnetars with PINNs

Key points

- PINNs are suitable for elliptic problems
- Relatively accurate and reliable
- Can be trained for arbitrary solutions → speed advantage
- Errors can be estimated through discretisation
- Connect regions with vastly different physical conditions
- New solutions for magnetothermal evolution
- Pave the road for extension to 3D



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Motivation

5 Pulsars with PINNs

Context

- Rotation is important
- Current-sheets
- \mathcal{T} must be consistent, not prescribed
- Additional constraints
- Axisymmetry

Pulsar equation

$$(1 - \beta^2) \Delta_{\text{Gs}} \mathcal{P} + 2\beta^2 q^2 \left(q \frac{\partial \mathcal{P}}{\partial q} + \mu \frac{\partial \mathcal{P}}{\partial \mu} \right) + \mathcal{T} \frac{d\mathcal{T}}{d\mathcal{P}} = 0$$

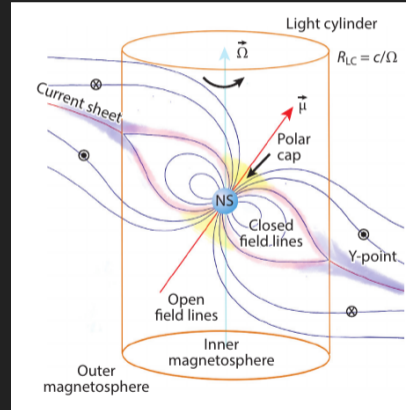


Figure 12: A sketch of a pulsar (Philippov and Kramer, 2022)



Results

5 Pulsars with PINNs

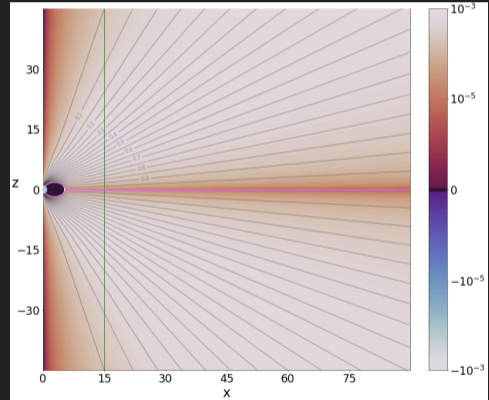
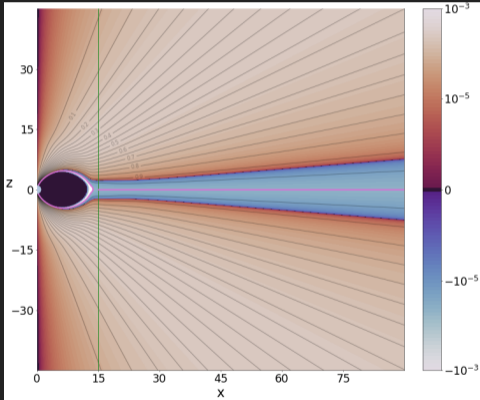


Figure 13: The classical pulsar magnetosphere
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Figure 14: A pulsar magnetosphere with the Y-point close to the star



Results

5 Pulsars with PINNs

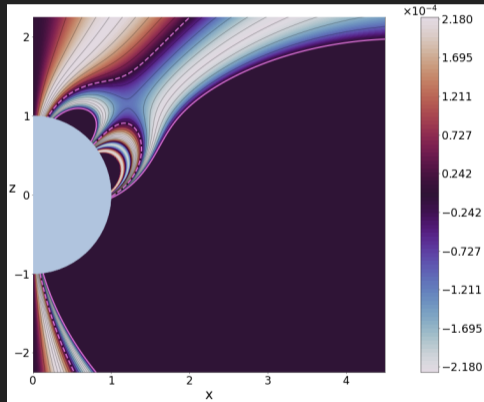
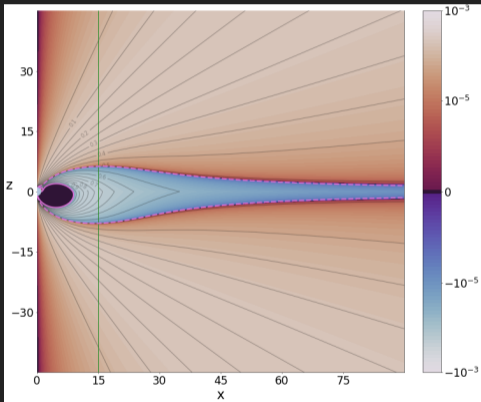


Figure 15: A multipolar pulsar magnetosphere

Figure 16: A close-up near the surface



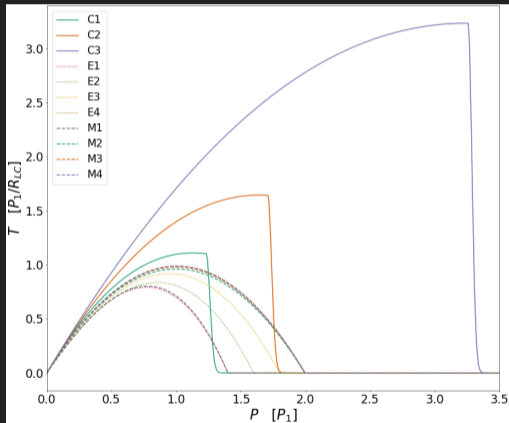
Discussion

5 Pulsars with PINNs



Discussion

5 Pulsars with PINNs



Models

- Excess energy $< 3\%$
- Models with fixed P_c have less energy
- Area represents luminosity
- Luminosity $< 3L_{dip}$
- For models with fixed P_c , not all Poynting flux escapes



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Ongoing and new projects

6 Future plans

PINNs

- Deeper understanding of how they work
- Improve PINN solvers

Magnetars

- Study stability of solutions
- Solve the 3D problem with PINNs
- Couple with 3D interior evolution

Pulsars

- Solve the full 3D problem

New projects

- Include magnetospheric radiation mechanisms
- Extend PINNs to other problems



Neutron star magnetospheres with classical and modern methods

Thank you



Publications

9 Publications




-  Stefanou, Petros, Jose A. Pons, and Pablo Cerdá-Durán (Feb. 2023). “Modelling 3D force-free neutron star magnetospheres”. In: *MNRAS* 518.4, pp. 6390–6400. DOI: [10.1093/mnras/stac3570](https://doi.org/10.1093/mnras/stac3570). arXiv: 2211.08957 [astro-ph.HE].
-  Urbán, Jorge F., Petros Stefanou, Clara Dehman, and José A. Pons (Sept. 2023). “Modelling force-free neutron star magnetospheres using physics-informed neural networks”. In: *MNRAS* 524.1, pp. 32–42. DOI: [10.1093/mnras/stad1810](https://doi.org/10.1093/mnras/stad1810). arXiv: 2303.11968 [astro-ph.HE].
-  Stefanou, Petros, Jorge F. Urbán, and José A. Pons (Nov. 2023). “Solving the pulsar equation using physics-informed neural networks”. In: *MNRAS* 526.1, pp. 1504–1511. DOI: [10.1093/mnras/stad2840](https://doi.org/10.1093/mnras/stad2840). arXiv: 2309.06410 [astro-ph.HE].



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10 Backup slides

► Backup slides

Optimisation process

- Trainable parameters Ω
- Loss function $\mathcal{J} = \|\mathbf{u}(\Omega) - \tilde{\mathbf{u}}\|$
- Ω are adjusted so that $\mathcal{J}(\Omega) \rightarrow 0$
- Corrections to Ω based on $\nabla_{\Omega} \mathcal{J}$
- Automatic differentiation

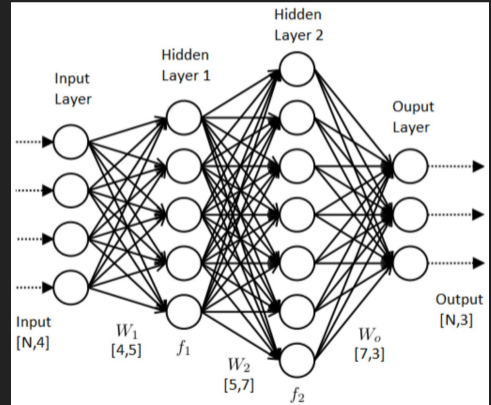


Figure 18: A neural network

Force-free case

- $\mathcal{T}(\mathcal{P}) = s_1 \mathcal{P} + s_2 \mathcal{P}^2$
- Up to seven multipoles $1 \leq l \leq 7$
- s_1, s_2 so that solutions exist
- Estimate error through discretisation

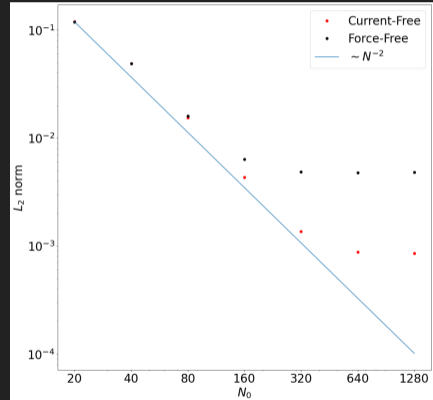


Figure 19: Discretisation error